

DARIUSZ PARTYKA (Lublin)

### Some Extremal Problems Concerning the Operator $B_\gamma$

**ABSTRACT.** Following [P2] we assign to each quasisymmetric automorphism  $\gamma$  of the unit circle  $\mathbf{T}$  a linear homeomorphic self-mapping  $B_\gamma$  of a Hilbert space  $(H, \|\cdot\|_H)$ . A complete solution to the following extremal problem is found: For which quasisymmetric automorphisms  $\gamma$  of  $\mathbf{T}$ ,  $\|B_\gamma(f)\|_H = \sqrt{K(\gamma)}$  for some  $f \in H$  with  $\|f\|_H = 1$ ? Here  $K(\gamma)$  stands for the maximal dilatation of an extremal quasiconformal extension of  $\gamma$  to the unit disk. As an application a relation between the Schober constant  $\lambda(\Gamma)$  of a quasicircle  $\Gamma \subset \mathbb{C}$  and an extremal quasiconformal reflection in  $\Gamma$  is established.

**0. Introduction.** Given a domain  $\Omega$  in the extended complex plain  $\hat{\mathbb{C}} := \mathbb{C} \cup \{\infty\}$  we denote by  $H(\Omega)$  ( $A(\Omega)$ ) the class of all complex-valued harmonic (analytic) functions on  $\Omega$ . If a function  $F : \Omega \rightarrow \mathbb{C}$  has partial derivatives for almost every (a.e. for short)  $z = x + iy \in \Omega$  then the Dirichlet integral  $\mathcal{D}_\Omega[F]$  of  $F$  is defined by

$$(0.1) \quad \mathcal{D}_\Omega[F] := \int_\Omega (|\partial_x F|^2 + |\partial_y F|^2) dS = 2 \int_\Omega (|\partial F|^2 + |\bar{\partial} F|^2) dS,$$

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where  $dS := dx dy$  and

$$(0.2) \quad \begin{aligned} \partial_x F &:= \frac{\partial F}{\partial x}, \quad \partial_y F := \frac{\partial F}{\partial y}, \\ \partial F &:= \frac{1}{2}(\partial_x F - i\partial_y F), \quad \bar{\partial} F := \frac{1}{2}(\partial_x F + i\partial_y F). \end{aligned}$$

The class  $\dot{A}^2(\Omega) := \{F \in A(\Omega) : \mathcal{D}_\Omega[F] < \infty\}$  is a closed subspace of the space  $\dot{H}^2(\Omega) := \{F \in H(\Omega) : \mathcal{D}_\Omega[F] < \infty\}$  in the pseudo-norm  $\|F\|_X := \sqrt{\frac{1}{2}\mathcal{D}_\Omega[F]}$ ,  $F \in X := \dot{H}^2(\Omega)$ .

Suppose  $\Omega$  is bounded by a Jordan curve  $\Gamma = \partial\Omega$ . Given a function  $F : \Omega \rightarrow \mathbb{C}$  we define for every  $z \in \Gamma$ ,  $\hat{\partial}F(z) := \lim_{\Omega \ni u \rightarrow z} F(u)$  provided the limit exists, while  $\hat{\partial}F(z) := 0$  otherwise. Write  $\Delta := \{z \in \mathbb{C} : |z| < 1\}$  and  $\mathbf{T} := \{z \in \mathbb{C} : |z| = 1\}$  for the unit disk and its boundary, respectively.

In case  $\Omega = \Delta$  we will use the symbol  $\hat{\partial}_r F$  for the radial limiting values function of  $F$ , i.e. for every  $z \in \mathbf{T}$ ,  $\hat{\partial}_r F_r(z) := \lim_{t \rightarrow 1^-} F(tz)$  if the limit exists, while  $\hat{\partial}_r F_r(z) := 0$  otherwise.

Given  $K \geq 1$  we denote by  $\mathbb{Q}_\Delta(K)$  the class of all  $K$ -quasiconformal ( $K$ -qc. for brevity) self-mappings of  $\Delta$ , and let  $\mathbb{Q}_\Delta := \bigcup_{K \geq 1} \mathbb{Q}_\Delta(K)$ . It is well known that every  $\varphi \in \mathbb{Q}_\Delta$  has a continuous extension to  $\mathbf{T}$  and  $\hat{\partial}\varphi$  is a sense-preserving homeomorphic self-mapping of  $\mathbf{T}$ ; cf. [LV, p. 42]. Due to Krzyż the class  $\mathbb{Q}_\mathbf{T} := \{\hat{\partial}\varphi : \varphi \in \mathbb{Q}_\Delta\}$  has a very simple characterization by means of quasisymmetric automorphisms of  $\mathbf{T}$ ; cf. [K1] and [K2].

Another interesting characterization of the class  $\mathbb{Q}_\mathbf{T}$  by quasihomographies was introduced by Zajac; cf. [Z], see also [K3]. For  $K \geq 1$ , define  $\mathbb{Q}_\mathbf{T}(K) := \{\hat{\partial}\varphi : \varphi \in \mathbb{Q}_\Delta(K)\}$ . Thus  $\mathbb{Q}_\mathbf{T}(K)$  is the class of all quasisymmetric automorphisms of  $\mathbf{T}$  which admit a  $K$ -qc. extension to  $\Delta$ . The functional  $K[\varphi] := \inf\{K \geq 1 : \varphi \in \mathbb{Q}_\Delta(K)\}$  is the maximal dilatation of  $\varphi$ .

Analogously, for  $\gamma \in \mathbb{Q}_\mathbf{T}$  we set  $K(\gamma) := \inf\{K \geq 1 : \gamma \in \mathbb{Q}_\mathbf{T}(K)\}$ . In both definitions  $\inf$  may be replaced by  $\min$  because of the compactness of the class  $\{\varphi \in \mathbb{Q}_\Delta(K) : \hat{\partial}\varphi = \gamma\}$  in the uniform convergence topology on  $\Delta$ ; cf. [LV, p. 73]. Thus  $K(\gamma)$  is the maximal dilatation of an extremal qc. extension  $\varphi$  of  $\gamma \in \mathbb{Q}_\mathbf{T}$  to  $\Delta$ ; extremal means that  $\varphi \in \mathbb{Q}_\Delta(K[\gamma])$ . For  $p \geq 1$ , we adopt the usual notation  $L^p(\mathbf{T})$  for the class of all functions  $f : \mathbf{T} \rightarrow \mathbb{C}$ ,  $p$ -integrable on  $\Gamma$  with respect to the Lebesgue arc-length measure, i.e.  $\|f\|_p := (\int_\Gamma |f(z)|^p |dz|)^{1/p} < \infty$ .

The notation  $f \doteq g$ ,  $f, g \in L^1(\mathbf{T})$ , means that  $f - g$  equals a constant almost everywhere (a.e. for brevity) on  $\mathbf{T}$ . It is clear that  $\doteq$  is an equivalence relation in  $L^1(\mathbf{T})$ , and let  $L^1(\mathbf{T}) := \{[f/\doteq] : f \in L^1(\mathbf{T})\}$  stand for the quotient space  $L^1(\mathbf{T})/\doteq$ . Recall that for every  $f \in L^1(\mathbf{T})$  and  $z \in \Delta$  the

Schwarz and Poisson formulas read as

$$(0.3) \quad f_{\Delta}(z) := \frac{1}{2\pi} \int_{\mathbf{T}} f(u) \frac{u+z}{u-z} |du| = a_0(f) + \sum_{n=1}^{\infty} a_n(f) z^n,$$

$$(0.4) \quad \mathcal{P}[f](z) := \frac{1}{2\pi} \int_{\mathbf{T}} f(u) \operatorname{Re} \frac{u+z}{u-z} |du|,$$

where

$$(0.5) \quad \begin{aligned} a_0(f) &:= \frac{1}{2\pi} \int_{\mathbf{T}} f(u) |du|, \\ a_n(f) &:= \frac{1}{\pi} \int_{\mathbf{T}} \bar{u}^n f(u) |du|, \quad n = 1, 2, \dots \end{aligned}$$

Obviously, for every  $f \in L^1(\mathbf{T})$ ,  $f_{\Delta} \in A(\Delta)$  and  $\mathcal{P}[f] \in H(\Delta)$ . According to the Poisson integral properties, for each  $f \in L^1(\mathbf{T})$  we have  $\hat{\partial}_r \mathcal{P}[f] = f$  a.e. on  $\mathbf{T}$ ; cf. [Du, p. 5], [R, Sect. 11.12]. Therefore, the operator  $\mathcal{S}_0 : L^1(\mathbf{T}) \rightarrow L^1(\mathbf{T})$ ,

$$\mathcal{S}_0([f/\doteq]) := \hat{\partial}_r \mathcal{P}[f] - \mathcal{P}[f](0), \quad f \in L^1(\mathbf{T}),$$

is a selector on the quotient space  $L^1(\mathbf{T})$ . We call it the Poisson selector.

Consider the class  $\dot{H}^2(\partial\Delta) := \{f \in L^1(\mathbf{T}) : \mathcal{P}[f] \in \dot{H}^2(\Delta)\}$ , and define the quotient space

$$\mathbf{H} := \operatorname{Re} \dot{H}^2(\partial\Delta) / \doteq = \{\mathbf{f} \in L^1(\mathbf{T}) : \mathcal{S}_0(\mathbf{f}) \in \operatorname{Re} \dot{H}^2(\partial\Delta)\}.$$

Here and subsequently,  $\operatorname{Re} X := \{\operatorname{Re} f : f \in X\}$  for any space  $X$  of complex-valued functions. If  $f \in \operatorname{Re} \dot{H}^2(\partial\Delta)$  then, by (0.3) and (0.5), we get

$$\begin{aligned} \|f\|_2^2 &= 2\pi |f_{\Delta}(0)|^2 + \pi \sum_{n=1}^{\infty} |a_n(f)|^2 \leq 2\pi |a_0(f)|^2 + \pi \sum_{n=1}^{\infty} n |a_n(f)|^2 \\ &= 2\pi |a_0(f)|^2 + \int_{\Delta} |(f_{\Delta})'|^2 dS < \infty, \end{aligned}$$

so that  $f \in L^2(\mathbf{T})$ . Therefore,  $(\mathbf{H}, \|\cdot\|_{\mathbf{H}})$  is a real Hilbert space, where

$$(0.6) \quad 2\|\mathbf{f}\|_{\mathbf{H}}^2 := \mathcal{D}[\mathcal{P}[\mathcal{S}_0(\mathbf{f})]] = \int_{\Delta} |(\mathcal{S}_0(\mathbf{f})_{\Delta})'|^2 dS.$$

For brevity we shall write  $\mathcal{D}[F]$  for the Dirichlet integral  $\mathcal{D}_{\Delta}[F]$ . We denote by  $\mathbb{P}$  the set of all complex polynomials. For a non-empty set  $K \subset \mathbb{C}$ , let  $\mathbb{P}(K) := \{P|_K : P \in \mathbb{P}\}$ . From (0.3), (0.5) and (0.6) it follows, in the standard way, that

$$(0.7) \quad \mathcal{S}_0(\mathbf{H}) \subset \operatorname{Re} L^2(\mathbf{T});$$

$$(0.8) \quad \|\mathcal{S}_0(\mathbf{f})\|_2 \leq \sqrt{2} \|\mathbf{f}\|_{\mathbf{H}}, \quad \mathbf{f} \in \mathbf{H};$$

$$(0.9) \quad \{\mathbf{f} : \mathcal{S}_0(\mathbf{f}) \in \operatorname{Re} \mathbb{P}(\mathbf{T})\} \text{ is a dense subspace of } \mathbf{H};$$

cf. [P6, Thm. 2.4.8] and [P5, Thm. 1.2]. Moreover, we can show that for every  $F \in \dot{A}^2(\Delta)$ ,  $F$  belongs to the Hardy class  $H^2$ , and so

$$(0.10) \quad F = (\operatorname{Re} \hat{\partial}_r F)_\Delta + i \operatorname{Im} F(0) \text{ and } \operatorname{Re} \hat{\partial}_r F \in \operatorname{Re} \dot{H}^2(\partial\Delta);$$

$$(0.11) \quad 2\|[\operatorname{Re} \hat{\partial}_r F / \doteq]\|_{\mathbf{H}}^2 = \mathcal{D}[\mathcal{P}[\operatorname{Re} \hat{\partial}_r F]] = \int_{\Delta} |((\operatorname{Re} \hat{\partial} F)_\Delta)'|^2 dS;$$

cf. [P6, Thm. 2.4.4]. We adopt the usual notation  $C(K)$  for the class of all complex-valued continuous functions on a set  $K \neq \emptyset$ . From Lemma 1.1 and (0.9) it follows that there exists a unique linear bounded operator  $B_\gamma : \mathbf{H} \rightarrow \mathbf{H}$  satisfying

$$(0.12) \quad B_\gamma([f / \doteq]) = [f \circ \gamma / \doteq], \quad f \in \operatorname{Re} C(\mathbf{T}) \cap \dot{H}^2(\partial\Delta).$$

Let  $\|T\|$  stand for the supremum norm of a linear operator  $T : \mathbf{H} \rightarrow \mathbf{H}$ , i.e.  $\|T\| := \sup\{\|T(f)\|_{\mathbf{H}} : f \in \mathbf{H} \text{ and } \|f\|_{\mathbf{H}} \leq 1\}$ . Following [BS] we will use the notation  $\tilde{\varphi}$  for the inverse mapping to a complex-valued mapping  $\varphi$  if it exists. By definition and by Theorem 1.2 we easily find that for any  $\gamma, \sigma \in \mathbb{Q}_{\mathbf{T}}$

$$(0.13) \quad \|B_\gamma\| \leq \sqrt{K(\gamma)};$$

$$(0.14) \quad B_{\gamma\sigma} = B_\sigma B_\gamma;$$

$$(0.15) \quad B_\gamma = B_\gamma^{-1};$$

$$(0.16) \quad B_{\operatorname{id}_{\mathbf{T}}} = I,$$

where  $\operatorname{id}_{\mathbf{T}} : \mathbf{T} \rightarrow \mathbf{T}$  and  $I : \mathbf{H} \rightarrow \mathbf{H}$  are identity mappings; cf. [P6, Corollary 2.5.4] and [P2, Lemma 1.1]. The properties (0.13) and (0.15) say that the operator  $B_\gamma$  is a linear homeomorphism of  $\mathbf{H}$  onto itself. Moreover, it turns out that

$$B_\gamma(f) = [S_0(f) \circ \gamma / \doteq], \quad f \in \mathbf{H}, \gamma \in \mathbb{Q}_{\mathbf{T}};$$

cf. [P6, formula (2.5.8)]. However, we will not use this fact in the sequel. In what follows we list four natural questions involving the supremum norm of the operator  $B_\gamma$ .

**Question 0.1.** For which  $\gamma \in \mathbb{Q}_{\mathbf{T}}$ ,  $\|B_\gamma\| = \sqrt{K(\gamma)}$ ?

**Question 0.2.** For which  $\gamma \in \mathbb{Q}_{\mathbf{T}}$ , does there exist  $f \in \mathbf{H}$  with  $\|f\|_{\mathbf{H}} = 1$  such that  $\|B_\gamma(f)\| = \|B_\gamma\|$ ? This question may be formulated equivalently: When  $\|B_\gamma\| = \max\{\|B_\gamma(f)\|_{\mathbf{H}} : f \in \mathbf{H} \text{ and } \|f\|_{\mathbf{H}} \leq 1\}$ ?

**Question 0.3.** For which  $\gamma \in \mathbb{Q}_{\mathbf{T}}$ ,  $\|B_\gamma(f)\|_{\mathbf{H}} = \sqrt{K(\gamma)}$  for some  $f \in \mathbf{H}$  with  $\|f\|_{\mathbf{H}} = 1$ ?

**Question 0.4.** Does there exist a constant  $c > 0$  such that for every  $\gamma \in \mathbb{Q}_{\mathbf{T}}$ ,  $\|B_\gamma\| - 1 \geq c(\sqrt{K(\gamma)} - 1)$ ?

In the next section we give a complete answer to the Question 0.3. In Section 2 we show that for some  $\gamma \in \mathbb{Q}_{\mathbf{T}}$ ,  $\|B_\gamma\|$  may be expressed by

the smallest positive eigenvalue  $\lambda_*(\gamma)$  of  $\gamma$ . The results obtained there are helpful in the next section. It turns out that the supremum norms  $\|\mathbf{B}_\gamma\|$  and  $\|\mathbf{B}_\gamma^{-1}\|$  are related to the Schober constant  $\lambda(\Gamma)$  of a certain quasicircle  $\Gamma \subset \mathbb{C}$  whose welding homeomorphism is  $\gamma \in \mathbb{Q}_T$ ; cf. Lemma 3.1. Thus the study of the Schober constant  $\lambda(\Gamma)$  can be reduced to the study of norms  $\|\mathbf{B}_\gamma\|$  and  $\|\mathbf{B}_\gamma^{-1}\|$ , which seems to be easier in some cases. As applications we present a few results in Section 3. The norm  $\|\mathbf{B}_\gamma\|$  is also closely related to the Grunsky-Kühnau constant  $\kappa$  (cf. [Kül, p. 383]) for a respective Grunsky matrix associated with  $\gamma$ . However, this topic will be studied in a forthcoming publication. This justifies studying the norm  $\|\mathbf{B}_\gamma\|$ . In the last section we give some comments to our subject.

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**1. The main result.** It is easily verified that

$$|a_1 b_1 + a_2 b_2|^2 + |a_1 \bar{b}_2 + a_2 \bar{b}_1|^2 \leq (|a_1|^2 + |a_2|^2)(|b_1| + |b_2|)^2$$

for any  $a_1, a_2, b_1, b_2 \in \mathbb{C}$ . The change of variables formula now shows that for all  $F \in \dot{H}^2(\Delta)$ ,  $K \geq 1$  and  $\varphi \in \mathbb{Q}_K(\Delta)$

$$\begin{aligned} (1.1) \quad \mathcal{D}_\Delta[F \circ \varphi] &= 2 \int_\Delta (|\partial(F \circ \varphi)|^2 + |\bar{\partial}(F \circ \varphi)|^2) dS \\ &\leq 2 \int_\Delta (|\partial F \circ \varphi|^2 + |\bar{\partial} F \circ \varphi|^2)(|\partial \varphi| + |\bar{\partial} \varphi|)^2 dS \\ &\leq 2K[\varphi] \int_\Delta (|\partial F \circ \varphi|^2 + |\bar{\partial} F \circ \varphi|^2)(|\partial \varphi|^2 - |\bar{\partial} \varphi|^2) dS \\ &= 2K[\varphi] \int_\Delta (|\partial F|^2 + |\bar{\partial} F|^2) dS = K[\varphi] \mathcal{D}_\Delta[F]. \end{aligned}$$

This means that the Dirichlet integral is quasi-invariant; cf. e.g. [A1, p. 18].

**Lemma 1.1.** *Given  $K \geq 1$  assume that  $\varphi \in \mathbb{Q}_K(\Delta)$ . Then for all functions  $F \in \text{Re } \mathbb{P}$  and  $P \in \text{Re } \dot{H}^2(\Delta)$ ,  $G := \mathcal{P}[\bar{\partial}(F \circ \varphi)] \in \text{Re } \dot{H}^2(\Delta)$  and*

$$(1.2) \quad \mathcal{D}[F \circ \varphi - G + P] = \mathcal{D}[F \circ \varphi - G] + \mathcal{D}[P].$$

*In particular,*

$$(1.3) \quad \mathcal{D}[G] = \mathcal{D}[F \circ \varphi] - \mathcal{D}[F \circ \varphi - G] \leq K \mathcal{D}[F].$$

**Proof.** Suppose  $K$ ,  $\varphi$ ,  $F$  and  $P$  satisfy the assumptions of our lemma and set  $\gamma := \bar{\partial} \varphi$ . The proof will be divided into two parts.

**Part I.** We first prove the lemma under the assumption that  $G \in \dot{H}^2(\Delta)$ . Since the class  $\text{ReP}(\Delta)$  is dense in  $\text{Re } \dot{H}^2(\Delta)$ , there exists a sequence  $P_n \in \text{ReP}$ ,  $n \in \mathbb{N}$ , such that

$$(1.4) \quad \mathcal{D}[P_n - P] \rightarrow 0, \quad n \rightarrow \infty.$$

For  $z \in \hat{\mathbb{C}}$  define  $\tilde{\varphi}(z) := \varphi(z)$  if  $z \in \Delta$ ,  $\tilde{\varphi}(z) := \gamma(z)$  if  $z \in \mathbf{T}$  and  $\tilde{\varphi}(z) := 1/\overline{\varphi(1/\bar{z})}$  if  $z \in \hat{\mathbb{C}} \setminus \overline{\Delta}$ . By the reflection principle for qc. mappings (see for instance [LV, p. 47]),  $\tilde{\varphi}$  is a qc. self-mapping of the extended complex plane  $\hat{\mathbb{C}}$ . For every  $t \in \mathbb{R}$  the set

$$\ell_\xi(t) := \{z \in \mathbb{C} : \text{Re}(z - t\xi)\bar{\xi} = 0\}, \quad \xi \in \mathbb{C} \setminus \{0\},$$

is the straight line passing through the point  $t\xi$  and orthogonal to the straight line  $\{s\xi : s \in \mathbb{R}\}$ . Since  $\tilde{\varphi}$  has the ACL-property (for the definition cf. e.g. [LV, p. 127 and 162]), it is absolutely continuous on almost every chord parallel to either of the coordinate axes, i.e.  $\tilde{\varphi}$  is absolutely continuous on  $\ell_\xi(t) \cap \Delta$  for a.e.  $t \in [-1, 1]$ ,  $\xi = 1, i$ . By definition,  $\hat{\partial}G = F \circ \gamma$ , and so  $\hat{\partial}(F \circ \varphi - G) = 0$  on  $\mathbf{T}$ . Moreover, by our assumption,  $\mathcal{D}[G] < \infty$ , so that for almost every  $y \in [-1, 1]$  and  $x \in [-1, 1]$

$$\int_{\Delta \cap \ell_i(y)} |\partial_x(F \circ \varphi - G)| < \infty \quad \text{and} \quad \int_{\Delta \cap \ell_1(x)} |\partial_y(F \circ \varphi - G)| < \infty.$$

Fix  $n \in \mathbb{N}$ . We may now integrate by parts to conclude that for a. e.  $y \in [-1, 1]$

$$\int_{\Delta \cap \ell_i(y)} \partial_x(F \circ \varphi - G) \partial_x P_n dx = - \int_{\Delta \cap \ell_i(y)} (F \circ \varphi - G) \partial_{xx}^2 P_n dx$$

and for a.e.  $x \in [-1, 1]$

$$\int_{\Delta \cap \ell_1(x)} \partial_y(F \circ \varphi - G) \partial_y P_n dx = - \int_{\Delta \cap \ell_1(x)} (F \circ \varphi - G) \partial_{yy}^2 P_n dx,$$

where  $\partial_{xx}^2 := \partial_x \partial_x$  and  $\partial_{yy}^2 := \partial_y \partial_y$ . Fubini's theorem then implies

$$\begin{aligned} & \int_{\Delta} (\partial_x(F \circ \varphi - G) \partial_x P_n + \partial_y(F \circ \varphi - G) \partial_y P_n) dS \\ &= \int_{-1}^1 \left( \int_{\Delta \cap \ell_1(y)} \partial_x(F \circ \varphi - G) \partial_x P_n dx \right) dy \\ &+ \int_{-1}^1 \left( \int_{\Delta \cap \ell_1(x)} \partial_y(F \circ \varphi - G) \partial_y P_n dy \right) dx \\ &= - \int_{-1}^1 \left( \int_{\Delta \cap \ell_1(y)} (F \circ \varphi - G) \partial_{xx}^2 P_n dx \right) dy \\ &- \int_{-1}^1 \left( \int_{\Delta \cap \ell_1(x)} (F \circ \varphi - G) \partial_{yy}^2 P_n dy \right) dx \\ &= - \int_{\Delta} (F \circ \varphi - G) (\partial_{xx}^2 P_n + \partial_{yy}^2 P_n) dS = 0, \end{aligned}$$

because  $P_n$  is a harmonic function on  $\Delta$ . Hence

$$\begin{aligned} \mathcal{D}[F \circ \varphi - G + P_n] &= \mathcal{D}[F \circ \varphi - G] + \mathcal{D}[P_n] \\ &+ 2 \int_{\Delta} (\partial_x(F \circ \varphi - G) \partial_x P_n + \partial_y(F \circ \varphi - G) \partial_y P_n) dS \\ &= \mathcal{D}[F \circ \varphi - G] + \mathcal{D}[P_n]. \end{aligned}$$

A passage to the limit now implies, by (1.4), that

$$\begin{aligned} \mathcal{D}[F \circ \varphi - G + P] &= \lim_{n \rightarrow \infty} \mathcal{D}[F \circ \varphi - G + P_n] \\ &= \mathcal{D}[F \circ \varphi - G] + \lim_{n \rightarrow \infty} \mathcal{D}[P_n] = \mathcal{D}[F \circ \varphi - G] + \mathcal{D}[P], \end{aligned}$$

and this is precisely the equality (1.2). Setting  $P := G$  in (1.2) we obtain the equality in (1.3). The inequality in (1.3) follows from (1.1).

**Part II.** We complete the proof by showing the first part of our assertion, i.e. we prove that  $G$  always belongs to  $\dot{H}^2(\Delta)$ . Let  $\mathbf{A}_{\mathbf{T}}$  be the class of all homeomorphisms  $\sigma : \mathbf{T} \rightarrow \mathbf{T}$  which have a conformal extension to some open annulus containing  $\mathbf{T}$ . It is easy to check that each  $\sigma \in \mathbf{A}_{\mathbf{T}}$  is a quasimetric automorphism of  $\mathbf{T}$ , so that  $\mathbf{A}_{\mathbf{T}} \subset \mathbf{Q}_{\mathbf{T}}$ . The inclusion follows immediately also from the Fehlmann characterization of the class  $\mathbf{Q}_{\mathbf{T}}$ ; cf. [F1, Thm. 3.1] and [F2]. It turns out that there exist a constant  $K^* \geq 1$  and a sequence  $\gamma_n \in \mathbf{A}_{\mathbf{T}} \cap \mathbf{Q}_{\mathbf{T}}(K^*)$ ,  $n \in \mathbf{N}$ , satisfying

$$(1.5) \quad \lim_{n \rightarrow \infty} \gamma_n(z) = \gamma(z), \quad z \in \mathbf{T};$$

cf. [P6, Lemma 3.1.3] and [P5, Thm. 2.1]. For  $n \in \mathbb{N}$  define  $G_n := \mathcal{P}[F \circ \gamma_n]$ . Fix  $n \in \mathbb{N}$ . It is easily seen from the Douglas formula that  $\mathcal{D}[G_n] < \infty$ ; cf. [D] and [A2, Thm. 2-5, p. 32]. However, this can be obtained in a more direct way as below. Integrating by parts we have for each  $k \in \mathbb{N}$

$$\begin{aligned} a_k &:= \frac{1}{\pi} \int_{\mathbb{T}} F \circ \gamma_n(u) \bar{u}^k |du| = \frac{1}{\pi} \int_0^{2\pi} F \circ \gamma_n(e^{it}) e^{-ikt} dt \\ &= -\frac{1}{\pi k^2} \int_0^{2\pi} e^{-ikt} \frac{d^2}{dt^2} F \circ \gamma_n(e^{it}) dt. \end{aligned}$$

Hence

$$\begin{aligned} \mathcal{D}[G_n] &= \int_{\Delta} |((F \circ \gamma_n)_{\Delta})'|^2 dS = \pi \sum_{k=1}^{\infty} k |a_k|^2 \\ &\leq 4\pi \left( \max_{0 \leq t \leq 2\pi} \left| \frac{d^2}{dt^2} F \circ \gamma_n(e^{it}) \right| \right)^2 \sum_{k=1}^{\infty} \frac{1}{k^3} < \infty, \end{aligned}$$

and we can use Part I to obtain

$$(1.6) \quad \mathcal{D}[G_n] \leq K^* \mathcal{D}[F].$$

From (1.5) it follows that  $\gamma_n$  is uniformly convergent to  $\gamma$ , and consequently for every  $z \in \Delta$ ,  $\lim_{n \rightarrow \infty} \partial G_n(z) = \partial G(z)$  and  $\lim_{n \rightarrow \infty} \bar{\partial} G_n(z) = \bar{\partial} G(z)$ . Then (1.6) shows, by Fatou's lemma, that  $\mathcal{D}[G] \leq \liminf_{n \rightarrow \infty} \mathcal{D}[G_n] \leq K^* \mathcal{D}[F] < \infty$ , which is our claim. Combining Parts I and II yields the assertion of the lemma.  $\square$

**Theorem 1.2.** *Given  $\gamma \in \mathbb{Q}_{\mathbb{T}}$  assume that  $\varphi$  is its qc. extension to  $\Delta$ . Then for every  $f \in H$*

$$(1.7) \quad 2\|B_{\gamma}(f)\|_H^2 = \mathcal{D}[F \circ \varphi] - \mathcal{D}[F \circ \varphi - G],$$

where  $F := \mathcal{P}[\mathcal{S}_0(f)]$  and  $G := \mathcal{P}[\mathcal{S}_0 B_{\gamma}(f)]$ . In particular,

$$(1.8) \quad \|B_{\gamma}(f)\|_H \leq \sqrt{K(\gamma)} \|f\|_H.$$

**Proof.** Suppose  $\gamma$  and  $\varphi$  are as in the assumption and fix  $f \in H$ . By (0.9) there exists a sequence  $f_n := [f_n / \dagger]$ ,  $f_n \in \text{Re } \mathbb{P}(\mathbb{T})$ ,  $n \in \mathbb{N}$ , such that

$$(1.9) \quad \|f - f_n\|_H \rightarrow 0 \quad \text{as } n \rightarrow \infty.$$



Then, by continuity,

$$(1.10) \quad \|B_\gamma(\mathbf{f}) - B_\gamma(\mathbf{f}_n)\|_H \rightarrow 0 \quad \text{as } n \rightarrow \infty .$$

For  $n \in \mathbb{N}$  set  $F_n := \mathcal{P}[f_n]$  and  $G_n := \mathcal{P}[f_n \circ \gamma]$ . From (1.10) and (0.12) it follows that

$$(1.11) \quad 2\|B_\gamma(\mathbf{f})\|_H^2 = 2 \lim_{n \rightarrow \infty} \|B_\gamma(\mathbf{f}_n)\|_H^2 = \lim_{n \rightarrow \infty} \mathcal{D}[G_n] .$$

By (1.1), (1.9), (0.6) and the Minkowski inequality we have

$$\begin{aligned} |\mathcal{D}[F \circ \varphi]^{1/2} - \mathcal{D}[F_n \circ \varphi]^{1/2}| &\leq \mathcal{D}[F \circ \varphi - F_n \circ \varphi]^{1/2} = \mathcal{D}[(F - F_n) \circ \varphi]^{1/2} \\ &\leq \sqrt{K[\varphi]} \mathcal{D}[(F - F_n)]^{1/2} = \sqrt{2K[\varphi]} \|\mathbf{f} - \mathbf{f}_n\|_H \rightarrow 0 \quad \text{as } n \rightarrow \infty , \end{aligned}$$

and consequently

$$(1.12) \quad \lim_{n \rightarrow \infty} \mathcal{D}[F_n \circ \varphi] = \mathcal{D}[F \circ \varphi] .$$

In the similar way we show that

$$\begin{aligned} |\mathcal{D}[F \circ \varphi - G]^{1/2} - \mathcal{D}[F_n \circ \varphi - G_n]^{1/2}| &\leq \mathcal{D}[F \circ \varphi - F_n \circ \varphi + G_n - G]^{1/2} \\ &\leq \mathcal{D}[(F - F_n) \circ \varphi]^{1/2} + \mathcal{D}[G - G_n]^{1/2} \\ &\leq \sqrt{2K[\varphi]} \|\mathbf{f} - \mathbf{f}_n\|_H + \sqrt{2} \|B_\gamma(\mathbf{f}) - B_\gamma(\mathbf{f}_n)\|_H \rightarrow 0 \quad \text{as } n \rightarrow \infty , \end{aligned}$$

and so

$$(1.13) \quad \lim_{n \rightarrow \infty} \mathcal{D}[F_n \circ \varphi - G_n] = \mathcal{D}[F \circ \varphi - G] .$$

From Lemma 1.1 we conclude that for every  $n \in \mathbb{N}$ ,  $\mathcal{D}[G_n] = \mathcal{D}[F_n \circ \varphi] - \mathcal{D}[F_n \circ \varphi - G_n]$ . Combining this with (1.11), (1.12) and (1.13) we obtain (1.7). It is a well known fact that for any  $K \geq 1$  the class  $\{\varphi \in \mathbb{Q}_\Delta(K) : \bar{\partial}\varphi = \gamma\}$  is compact in the uniform convergence topology on  $\Delta$ ; cf. e.g. [LV, p. 73]. Therefore there exists an extremal  $K(\gamma)$ -qc. extension  $\psi$  of  $\gamma$  to  $\Delta$ . Setting  $\varphi := \psi$  in (1.7) yields (1.8).  $\square$

We recall that a qc. self-mapping  $\psi$  of  $\Delta$  is said to be a regular Teichmüller mapping if there exists a non-zero function  $F \in A(\Delta)$  and a constant  $k$ ,  $0 \leq k < 1$ , such that the complex dilatation of  $\psi$  is of the form

$$(1.14) \quad \frac{\bar{\partial}\psi}{\partial\psi} = k \frac{\bar{F}}{|F|} \quad \text{a.e. on } \Delta .$$

We are now in a position to answer the Question 0.3.

**Theorem 1.3.** *Let  $\gamma \in \mathbb{Q}_{\mathbb{T}}$  and let  $0 \leq k < 1$ . If  $f \in H$  satisfies  $\|f\|_H > 0$  and if  $\gamma$  admits a regular qc. Teichmüller extension  $\psi$  to  $\Delta$  with the complex dilatation*

$$(1.15) \quad \frac{\bar{\partial}\psi}{\partial\psi} = k \frac{\overline{(\mathcal{S}_0(f)_\Delta)'}}{(\mathcal{S}_0(f)_\Delta)'} \quad \text{a.e. on } \Delta ,$$

*then there exists  $g \in H$  such that  $\|g\|_H = 1$  and*

$$(1.16) \quad \|B_\gamma(g)\|_H^2 = K(\gamma) = \frac{1+k}{1-k} .$$

*In particular,  $\|B_\gamma\| = \sqrt{K(\gamma)}$ .*

*Conversely, if  $g \in H$  and  $k$  satisfy  $\|g\|_H = 1$  and (1.16), then  $\gamma$  admits a regular qc. Teichmüller extension  $\psi$  to  $\Delta$  with the complex dilatation (1.15), where  $f := B_\gamma(g)$ . Moreover,  $\psi$  is uniquely extremal.*

**Proof.** Assume  $\gamma \in \mathbb{Q}_{\mathbb{T}}$  admits a qc. extension  $\psi$  to  $\Delta$  with the complex dilatation (1.15). Let  $F := \mathcal{S}_0(f)_\Delta$  and  $G : \Delta \rightarrow \mathbb{C}$  be a function satisfying the equation

$$(1.17) \quad (1 - k)G \circ \psi = k\bar{F} + F .$$

Differentiating both sides of this equality we get

$$\begin{aligned} (\partial G) \circ \psi \partial\psi + (\bar{\partial}G) \circ \psi \partial\bar{\psi} &= (1 - k)^{-1}F' , \\ (\partial G) \circ \psi \bar{\partial}\psi + (\bar{\partial}G) \circ \psi \bar{\partial}\bar{\psi} &= (1 - k)^{-1}k\bar{F}' . \end{aligned}$$

Since  $\partial\psi \bar{\partial}\bar{\psi} - \bar{\partial}\psi \partial\bar{\psi} = \partial\psi \bar{\partial}\bar{\psi} - \bar{\partial}\psi \partial\bar{\psi} = |\partial\psi|^2 - |\bar{\partial}\psi|^2 > 0$  a.e. on  $\Delta$ , (1.15) shows that  $\bar{\partial}G = 0$  a.e. on  $\Delta$ . In this way the function  $G$  is analytic on  $\Delta$ ; cf. [A1, p. 33]. Moreover, by (1.1) we have

$$\begin{aligned} 2(1 - k)^2 \int_{\Delta} |G'|^2 dS &= \mathcal{D}[(1 - k)G] = \mathcal{D}[(k\bar{F} + F) \circ \check{\psi}] \\ &\leq K[\check{\psi}]\mathcal{D}[k\bar{F} + F] = 2K[\psi](1 + k^2) \int_{\Delta} |F'|^2 dS < \infty . \end{aligned}$$

Thus  $G \in \dot{A}^2(\Delta)$  and, by the definition of  $H$ , there exists  $g \in H$  such that  $G - G(0) = \mathcal{S}_0(g)_\Delta$ . By (0.9) there exists a sequence  $g_n \in H$ ,  $n \in \mathbb{N}$ , such that  $g_n := \mathcal{S}_0(g_n) \in \text{Re}\mathbb{P}(\mathbb{T})$  and

$$(1.18) \quad \|g - g_n\|_H \rightarrow 0 \quad \text{as } n \rightarrow \infty .$$

From (1.17) we see that  $G \circ \psi \in \dot{H}^2(\Delta)$ , so that

$$(1.19) \quad \operatorname{Re} G \circ \psi = \mathcal{P}[\mathcal{S}_0(\mathbf{h})] + c$$

for some  $\mathbf{h} \in \mathbf{H}$  and  $c \in \mathbb{R}$ . Moreover, from Lemma 1.1, (1.1), (0.6) and (1.18) it follows that

$$(1.20) \quad \begin{aligned} 2\|\mathbf{h} - \mathbf{B}_\gamma(\mathbf{g}_n)\|_{\mathbf{H}}^2 &= \mathcal{D}[\operatorname{Re} G \circ \psi - \mathcal{P}[g_n \circ \gamma]] \leq \mathcal{D}[\operatorname{Re} G \circ \psi - \mathcal{P}[g_n] \circ \psi] \\ &\leq K[\psi] \mathcal{D}[\operatorname{Re} G - \mathcal{P}[g_n]] = 2K[\psi] \|\mathbf{g} - \mathbf{g}_n\|_{\mathbf{H}}^2 \rightarrow 0 \quad \text{as } n \rightarrow \infty. \end{aligned}$$

On the other hand, by the definition of the operator  $\mathbf{B}_\gamma$  and by (1.18), we obtain

$$\|\mathbf{B}_\gamma(\mathbf{g}_n) - \mathbf{B}_\gamma(\mathbf{g})\|_{\mathbf{H}} \rightarrow 0 \quad \text{as } n \rightarrow \infty.$$

Combining this with (1.20) we conclude that

$$(1.21) \quad \mathbf{B}_\gamma(\mathbf{g}) = \mathbf{h}.$$

Theorem 1.2 now shows, by (1.19), (1.21), (0.6) and (1.17), that

$$(1.22) \quad \begin{aligned} 2\|\mathbf{B}_\gamma(\mathbf{g})\|_{\mathbf{H}}^2 &= \mathcal{D}[\operatorname{Re} G \circ \psi] - \mathcal{D}[\operatorname{Re} G \circ \psi - \mathcal{P}[\mathcal{S}_0 \mathbf{B}_\gamma(\mathbf{g})]] \\ &= \mathcal{D}[\operatorname{Re} G \circ \psi] - \mathcal{D}[\operatorname{Re} G \circ \psi - \mathcal{P}[\mathcal{S}_0(\mathbf{h})]] = \mathcal{D}[\operatorname{Re} G \circ \psi] \\ &= \left(\frac{1+k}{1-k}\right)^2 \mathcal{D}[\operatorname{Re} F] = 2\left(\frac{1+k}{1-k}\right)^2 \|\mathbf{f}\|_{\mathbf{H}}^2. \end{aligned}$$

Since  $G$  is analytic, we see that

$$G' \circ \psi \partial \psi = (1-k)^{-1} F' \quad \text{and} \quad G' \circ \psi \bar{\partial} \psi = (1-k)^{-1} k \bar{F}',$$

and hence, by (0.6), that

$$(1.23) \quad \begin{aligned} 2\frac{1+k}{1-k} \|\mathbf{f}\|_{\mathbf{H}}^2 &= \frac{1+k}{1-k} \int_{\Delta} |F'|^2 dS \\ &= \frac{1}{(1-k)^2} \int_{\Delta} |F'|^2 dS - \frac{k^2}{(1-k)^2} \int_{\Delta} |\bar{F}'|^2 dS \\ &= \int_{\Delta} (|G' \circ \psi \partial \psi|^2 - |G' \circ \psi \bar{\partial} \psi|^2) dS \\ &= \int_{\Delta} |G' \circ \psi|^2 (|\partial \psi|^2 - |\bar{\partial} \psi|^2) dS = \int_{\Delta} |G'|^2 dS = 2\|\mathbf{g}\|_{\mathbf{H}}^2. \end{aligned}$$

Combining this with (1.8) and (1.22) we obtain

$$(1.24) \quad \|\mathbf{B}_\gamma(\mathbf{g})\|_{\mathbf{H}}^2 \leq K(\gamma) \|\mathbf{g}\|_{\mathbf{H}}^2 \leq \left(\frac{1+k}{1-k}\right)^2 \|\mathbf{f}\|_{\mathbf{H}}^2 = \|\mathbf{B}_\gamma(\mathbf{g})\|_{\mathbf{H}}^2.$$

From (1.23) and (1.24) it follows that

$$\|B_\gamma^\bullet(g)\|_H^2 = K(\gamma)\|g\|_H^2 = (1+k)(1-k)^{-1}\|g\|_H^2.$$

Replacing now  $g$  by  $g/\|g\|_H$  we obtain (1.16). That  $\|B_\gamma\| = \sqrt{K(\gamma)}$  follows from Theorem 1.2 and (1.16), which proves the first part of Theorem 1.3.

Suppose now that  $g \in H$  and  $k$  satisfy  $\|g\|_H = 1$  and (1.16). Since for any  $K \geq 1$  the class  $\{\varphi \in Q_\Delta(K) : \bar{\partial}\varphi = \gamma\}$  is compact in the uniform convergence topology on  $\Delta$  (cf. e.g. [LV, p. 73]), there exists an extremal  $K(\gamma)$ -qc. extension  $\psi$  of  $\gamma$  to  $\Delta$ . Set  $F := (S_0 B_\gamma(g))_\Delta$  and  $G := (S_0(g))_\Delta$ . Then Theorem 1.2 shows, by (1.1) and (0.6), that

$$\begin{aligned} \mathcal{D}[\operatorname{Re} G \circ \psi - \operatorname{Re} F] &= \mathcal{D}[\operatorname{Re} G \circ \psi] - 2\|B_\gamma(g)\|_H^2 \leq \\ K[\psi]\mathcal{D}[\operatorname{Re} G] - 2K(\gamma) &= 2K(\gamma)(\|g\|_H^2 - 1) = 0. \end{aligned}$$

Hence the equality  $\operatorname{Re} G \circ \psi = \operatorname{Re} F + c$  holds a.e. on  $\Delta$  for some constant  $c \in \mathbb{R}$ . Differentiating both sides of this equality we get

$$(1.25) \quad G' \circ \psi \partial\psi + \bar{G}' \circ \psi \bar{\partial}\bar{\psi} = F'.$$

Hence,

$$\begin{aligned} (1.26) \quad |F'|^2 &= |G' \circ \psi|^2 |\partial\psi|^2 + |\bar{G}' \circ \psi|^2 |\bar{\partial}\bar{\psi}|^2 + (G' \circ \psi)^2 \partial\psi \bar{\partial}\bar{\psi} + (\bar{G}' \circ \psi)^2 \bar{\partial}\bar{\psi} \partial\psi \\ &= |G' \circ \psi|^2 (|\partial\psi| + |\bar{\partial}\bar{\psi}|)^2 - Q = |G' \circ \psi|^2 \frac{|\partial\psi| + |\bar{\partial}\bar{\psi}|}{|\partial\psi| - |\bar{\partial}\bar{\psi}|} (|\partial\psi|^2 - |\bar{\partial}\bar{\psi}|^2) - Q \end{aligned}$$

a.e. on  $\Delta$ , where  $Q := 2|G' \circ \psi|^2 |\partial\psi| |\bar{\partial}\bar{\psi}| - (G' \circ \psi)^2 \partial\psi \bar{\partial}\bar{\psi} - (\bar{G}' \circ \psi)^2 \bar{\partial}\bar{\psi} \partial\psi$  a.e. on  $\Delta$ . Since  $\psi$  is  $K(\gamma)$ -qc.,  $(|\partial\psi| + |\bar{\partial}\bar{\psi}|)(|\partial\psi| - |\bar{\partial}\bar{\psi}|)^{-1} \leq K(\gamma)$  a.e. on  $\Delta$ . Combining this with (1.16) and (1.26) we obtain by (0.6)

$$\begin{aligned} 2K(\gamma) &= 2\|B_\gamma(g)\|_H^2 = \int_\Delta |F'|^2 dS \\ &\leq K(\gamma) \int_\Delta |G' \circ \psi|^2 (|\partial\psi|^2 - |\bar{\partial}\bar{\psi}|^2) dS - \int_\Delta Q dS \\ &= K(\gamma) \int_\Delta |G'|^2 dS - \int_\Delta Q dS = 2K(\gamma)\|g\|_H^2 - \int_\Delta Q dS \\ &= 2K(\gamma) - \int_\Delta Q dS \leq 2K(\gamma). \end{aligned}$$

The inequality is possible iff the equalities

$$|G' \circ \psi|^2 |\partial\psi| |\bar{\partial}\bar{\psi}| = (G' \circ \psi)^2 \partial\psi \bar{\partial}\bar{\psi} \quad \text{and} \quad \frac{|\bar{\partial}\bar{\psi}|}{|\partial\psi|} = k$$

hold a.e. on  $\Delta$ . Therefore,

$$(1.27) \quad \frac{\bar{\partial}\psi}{\partial\psi} = k \frac{\overline{G' \circ \psi \partial\psi}}{G' \circ \psi \partial\psi} \quad \text{a.e. on } \Delta .$$

Let  $\mathbf{f} := B_\gamma(\mathbf{g})$ . Then  $F = (\mathcal{S}_0(\mathbf{f}))_\Delta$ . We conclude from (1.27) and (1.25) that  $G' \circ \psi \partial\psi = (1 + k)^{-1} F'$ , hence that

$$(1.28) \quad \frac{\bar{\partial}\psi}{\partial\psi} = k \frac{\overline{F'}}{F'} \quad \text{a.e. on } \Delta ,$$

and finally that (1.15) holds. This ends the proof of the converse statement.

We now prove the uniqueness of the extremal extension  $\psi$ . Suppose  $\tilde{\psi}$  is another extremal  $K(\gamma)$ -qc. extension of  $\gamma$  to  $\Delta$ . Then the Beltrami equation (1.28) holds with  $\psi$  replaced by  $\tilde{\psi}$ . Hence  $\bar{\partial}(\tilde{\psi} \circ \psi) = 0$  a.e. on  $\Delta$ , and so  $\tilde{\psi} \circ \psi \in \mathbb{Q}_\Delta(1)$ . Since  $\hat{\partial}\psi = \hat{\partial}\tilde{\psi} = \gamma$ , we see that  $\psi = \tilde{\psi}$  on  $\Delta$ .  $\square$

**Corollary 1.4.** *If  $K \geq 1$  and if  $\mathbf{g} \in \mathbf{H}$  satisfies  $\|\mathbf{g}\|_{\mathbf{H}} = 1$ , then there exists  $\gamma \in \mathbb{Q}_{\mathbf{T}}$  such that*

$$(1.29) \quad \|\mathbf{B}_\gamma\|^2 = \|\mathbf{B}_\gamma(\mathbf{g})\|_{\mathbf{H}}^2 = K(\gamma) = K .$$

Moreover,  $\gamma$  admits a unique regular  $K(\gamma)$ -qc. Teichmüller extension  $\psi$  to  $\Delta$  with the complex dilatation given by (1.15), where  $k := (K - 1)/(K + 1)$  and  $\mathbf{f} := B_\gamma(\mathbf{g})$ .

**Proof.** Given  $K \geq 1$  and  $\mathbf{g} \in \mathbf{H}$ , let  $\|\mathbf{g}\|_{\mathbf{H}} = 1$  and let  $k$  and  $\mathbf{f}$  be as above. By the Mapping Theorem [LV, p. 194] (also cf. [B] and [LK, p. 45]), there exists a solution  $\psi$  of the Beltrami equation (1.15) being a  $K$ -qc. self-mapping of  $\Delta$ . Hence  $\gamma := \hat{\partial}\psi \in \mathbb{Q}_{\mathbf{T}}(K)$ . Theorem 1.3 now shows that (1.29) holds and  $\gamma$  admits a unique regular  $K(\gamma)$ -qc. Teichmüller extension  $\psi$  to  $\Delta$  with the complex dilatation given by (1.15).  $\square$

**2. The smallest eigenvalue of a quasymmetric automorphism of the unit circle.** If  $f \in \text{Re } L^1(\mathbf{T})$  then, by (0.3),  $\text{Im } f_\Delta$  is a real-valued harmonic function on  $\Delta$ . A classical result states that the function  $\text{Im } f_\Delta$  has a finite non-tangential limit a.e. on  $\mathbf{T}$  and

$$\hat{\partial}_r \text{Im } f_\Delta(z) = \lim_{r \rightarrow 1^-} \text{Im } f_\Delta(rz) = \frac{1}{\pi} \text{Re PV} \int_{\mathbf{T}} \frac{f(u)}{z - u} du$$

for a.e.  $z \in \mathbf{T}$ ; cf. e.g. [G, p. 103]. For every  $\mathbf{f} := [f/\div] \in \mathbf{H}$ , define

$$(2.1) \quad \mathbf{A}(\mathbf{f}) := [\hat{\partial}_r \text{Im } f_\Delta / \div] .$$

Since  $\mathcal{D}[\operatorname{Re} F] = \mathcal{D}[\operatorname{Im} F]$  for  $F \in \dot{A}^2(\Delta)$ , we conclude from (2.1) and the definition of the space  $\mathbf{H}$  that

$$(2.2) \quad \mathbf{A}(\mathbf{H}) = \mathbf{H} \quad , \quad \mathbf{A}^2 = -\mathbf{I} \quad \text{and} \quad \|\mathbf{A}\| = 1 \quad ,$$

and so the operator  $\mathbf{A}$  maps isometrically  $\mathbf{H}$  onto itself. Therefore, the operator

$$(2.3) \quad \mathbf{A}_\gamma := \mathbf{B}_\gamma \mathbf{A} \mathbf{B}_\gamma^{-1} \quad ,$$

called the generalized harmonic conjugation operator, is a linear homeomorphism of  $\mathbf{H}$  onto itself; cf. [P2]. We recall that a real number  $\lambda$  is said to be an eigenvalue of  $\gamma \in \mathbb{Q}_{\mathbf{T}}$  if there exists  $\mathbf{f} \in \mathbf{H}$  with  $\|\mathbf{f}\|_{\mathbf{H}} = 1$  such that

$$(2.4) \quad (\lambda + 1)\mathbf{A}(\mathbf{f}) = (\lambda - 1)\mathbf{A}_\gamma(\mathbf{f}) \quad ;$$

cf. [P3, Definition 1.1]. For every  $\gamma \in \mathbb{Q}_{\mathbf{T}}$  write  $\Lambda_\gamma^*$  for the set of all eigenvalues of  $\gamma$  and define

$$\lambda_*(\gamma) = \min\{\lambda > 0 : \lambda \in \Lambda_\gamma^*\}$$

whenever  $\Lambda_\gamma^* \neq \emptyset$  and the minimum exists, while  $\lambda_*(\gamma) = \infty$  otherwise. From [P3, Thm. 1.4] it follows that  $\lambda_*(\gamma) = \infty$  for  $\gamma \in \mathbb{Q}_{\mathbf{T}}(1)$ , and

$$(2.5) \quad \lambda_*(\gamma) \geq (K(\gamma) + 1)/(K(\gamma) - 1) \quad \text{for } \gamma \in \mathbb{Q}_{\mathbf{T}} \setminus \mathbb{Q}_{\mathbf{T}}(1) \quad .$$

A sufficient condition on  $\gamma$  for the equality in (2.5) to hold, was obtained in [P4, Thm. 2.2]. We use this result to show the following

**Theorem 2.1.** *Let  $\gamma \in \mathbb{Q}_{\mathbf{T}} \setminus \mathbb{Q}_{\mathbf{T}}(1)$ . Then*

$$(2.6) \quad \lambda_*(\gamma) = (K(\gamma) + 1)/(K(\gamma) - 1)$$

*iff there exists  $\mathbf{g} \in \mathbf{H}$  such that  $\|\mathbf{g}\|_{\mathbf{H}} = 1$  and that*

$$(2.7) \quad \|\mathbf{B}_\gamma(\mathbf{g})\|_{\mathbf{H}}^2 = K(\gamma) \quad .$$

**Proof.** Assume first that (2.7) holds. Then Theorem 1.3 shows that  $\gamma$  admits a regular qc. Teichmüller extension  $\psi$  to  $\Delta$  with the complex dilatation (1.15), where  $\mathbf{f} := \mathbf{B}_\gamma(\mathbf{g})$  and  $k := (K(\gamma) - 1)/(K(\gamma) + 1)$ . Therefore (2.6) follows from [P4, Thm. 2.2].

Conversely, assume that (2.6) holds. Then there exists  $\mathbf{f} \in \mathbf{H}$  such that  $\|\mathbf{f}\|_{\mathbf{H}} = 1$  and that (2.4) holds with  $\lambda$  replaced by  $\lambda_*(\gamma)$ . Hence, by (2.2) and (2.5) we have

$$(\lambda_*(\gamma) + 1)\|\mathbf{f}\|_{\mathbf{H}} = (\lambda_*(\gamma) - 1)\|\mathbf{A}_\gamma(\mathbf{f})\|_{\mathbf{H}}$$

and consequently, by (2.3) and (2.6),

$$K(\gamma) = \|B_\gamma A B_\gamma^{-1}(f)\|_H.$$

Set  $g := (\sqrt{K(\gamma)})^{-1} A B_\gamma^{-1}(f)$ . Theorem 1.2, (0.15) and (2.2) now imply that

$$\begin{aligned} \|g\|_H &\leq (\sqrt{K(\gamma)})^{-1} \|A\| \|B_\gamma^{-1}\| \|f\|_H = (\sqrt{K(\gamma)})^{-1} \|B_\gamma\| \\ &\leq (\sqrt{K(\gamma)})^{-1} \sqrt{K(\tilde{\gamma})} = 1; \end{aligned}$$

moreover,

$$\begin{aligned} K(\gamma) &= \|B_\gamma A B_\gamma^{-1}(f)\|_H = \sqrt{K(\gamma)} \|B_\gamma(g)\|_H \\ &\leq \sqrt{K(\gamma)} \|B_\gamma\| \|g\|_H \leq K(\gamma). \end{aligned}$$

Combining the above inequalities we see, by (0.13), that  $\|g\|_H = 1$  and the equality (2.7) holds.  $\square$

**Corollary 2.2.** *If  $\gamma \in \mathbb{Q}_T$  and if (2.7) holds for some  $g \in H$  such that  $\|g\|_H = 1$ , then there exists  $f \in H$  satisfying*

$$(2.8) \quad \|f\|_H = 1, \quad \|B_\gamma^{-1}(f)\|_H^2 = K(\gamma).$$

*In particular,  $\|B_\gamma\| = \|B_\gamma^{-1}\| = \sqrt{K(\gamma)}$ .*

**Proof.** Given  $\gamma \in \mathbb{Q}_T \setminus \mathbb{Q}_T(1)$  assume (2.7) holds for some  $g \in H$  such that  $\|g\|_H = 1$ . Therefore (2.6) holds by Theorem 2.1. From [P3, Thm. 1.4 (v)] it follows that  $\Lambda_\gamma^* = \Lambda_{\tilde{\gamma}}^*$ , and hence

$$\lambda_*(\gamma) = \min\{\lambda > 0 : \lambda \in \Lambda_\gamma^*\} = \min\{\lambda > 0 : \lambda \in \Lambda_{\tilde{\gamma}}^*\} = \lambda_*(\tilde{\gamma}).$$

Combining this with (2.6) we obtain

$$\lambda_*(\tilde{\gamma}) = \lambda_*(\gamma) = (K(\gamma) + 1)/(K(\gamma) - 1) = (K(\tilde{\gamma}) + 1)/(K(\tilde{\gamma}) - 1).$$

Applying Theorem 2.1 again, with  $\gamma$  replaced by  $\tilde{\gamma}$ , we see, by (0.15), that there exists  $f \in H$  satisfying (2.8). Moreover, combining (2.7) and (2.8) with (0.13) we have

$$\begin{aligned} \sqrt{K(\gamma)} &= \|B_\gamma(g)\|_H \leq \|B_\gamma\| \leq \sqrt{K(\gamma)}, \\ \sqrt{K(\tilde{\gamma})} &= \|B_\gamma(f)\|_H \leq \|B_\gamma\| \leq \sqrt{K(\tilde{\gamma})}, \end{aligned}$$

and hence  $\|B_\gamma\| = \sqrt{K(\gamma)} = \sqrt{K(\tilde{\gamma})} = \|B_\gamma\| = \|B_\gamma^{-1}\|$  as claimed.

If  $\gamma \in \mathbb{Q}_T(1)$  then, by (0.13) and (0.15), we obtain  $\|B_\gamma\| \leq \sqrt{K(\gamma)} = 1$  and  $\|B_\gamma^{-1}\| = \|B_\gamma\| \leq \sqrt{K(\tilde{\gamma})} = 1$ . Hence the operators  $B_\gamma$  and  $B_\gamma^{-1}$  are isometries of  $H$  onto itself, and the corollary follows.  $\square$

**3. The Schober constant  $\lambda_s(\Gamma)$ .** Given a Jordan curve  $\Gamma \subset \mathbb{C}$  write  $H(\Gamma)$  for the family of all real-valued functions  $F$  continuous on  $\hat{\mathbb{C}}$  and harmonic on  $\Omega \cup \Omega_* = \hat{\mathbb{C}} \setminus \Gamma$  which satisfy  $0 < \mathcal{D}_\Omega[F] + \mathcal{D}_{\Omega_*}[F] < \infty$ , where  $\Omega$  and  $\Omega_* \ni \infty$  are complementary domains to  $\Gamma$ . Define

$$\frac{1}{\lambda_s(\Gamma)} := \sup \left\{ \frac{|\mathcal{D}_\Omega[F] - \mathcal{D}_{\Omega_*}[F]|}{\mathcal{D}_\Omega[F] + \mathcal{D}_{\Omega_*}[F]} : F \in H(\Gamma) \right\}$$

provided the supremum is a positive number, while  $\lambda_s(\Gamma) = \infty$  otherwise. For a short survey of basic properties of the curve functional  $\lambda_s(\Gamma)$  we refer the reader to the Schober article [S]; also see the references given there. Let  $\Phi$  and  $\Phi_*$  denote conformal mappings of  $\Delta$  and  $\hat{\mathbb{C}} \setminus \bar{\Delta}$  onto the domains  $\Omega$  and  $\Omega_*$ , respectively. Such mappings exist by the Riemann mapping theorem; cf. for instance [R, Thm. 14.8]. Moreover, by the Taylor-Osgood-Carathéodory theorem both the mappings  $\Phi$  and  $\Phi_*$  have homeomorphic extensions to the closures  $\bar{\Omega}$  and  $\bar{\Omega}_*$ , respectively; cf. for instance [R, Thm. 14.19]. Then  $\gamma := \hat{\partial}\Phi_* \circ \hat{\partial}\Phi$  is a sense-preserving homeomorphic self-mapping of  $\mathbb{T}$ . We recall that every homeomorphism  $\gamma$  assigned to  $\Gamma$  in this way is said to be a welding homeomorphism of  $\Gamma \subset \mathbb{C}$ . The class of all welding homeomorphisms of  $\Gamma$  will be denoted by  $\text{Weld}(\Gamma)$ . For  $z \in \mathbb{C} \setminus \{0\}$  set  $h(z) := 1/z$ , and let  $h(0) := \infty$ ,  $h(\infty) := 0$ . If a Jordan curve  $\Gamma \subset \mathbb{C}$  admits a  $K$ -qc. reflection  $\Psi$  then  $\psi := \bar{h} \circ \check{\Phi}_* \circ \Psi \circ \Phi$  is a  $K$ -qc. extension of  $\gamma := \hat{\partial}\check{\Phi}_* \circ \hat{\partial}\Phi$  to  $\Delta$ . Conversely, if  $\psi$  is a  $K$ -qc. extension of  $\gamma$  to  $\Delta$  then  $\Psi$ ,

$$(3.1) \quad \Psi(z) := \begin{cases} \check{\Phi}_* \circ \bar{h} \circ \psi \circ \check{\Phi}(z) & , z \in \bar{\Omega} , \\ \check{\Psi}(z) & , z \in \Omega_* , \end{cases}$$

is a  $K$ -qc. reflection in  $\Gamma$ . Thus for every  $K \geq 1$ ,

$$(3.2) \quad \text{a Jordan curve } \Gamma \subset \mathbb{C} \text{ admits a } K\text{-qc. reflection iff } \text{Weld}(\Gamma) \subset \mathbb{Q}_{\mathbb{T}}(K) .$$

**Lemma 3.1.** *For every quasicircle  $\Gamma \subset \mathbb{C}$  the following equality holds*

$$(3.3) \quad \frac{1}{\lambda_s(\Gamma)} = \frac{(\max\{\|B_\gamma\|, \|B_\gamma^{-1}\|\})^2 - 1}{(\max\{\|B_\gamma\|, \|B_\gamma^{-1}\|\})^2 + 1} ,$$

where  $\gamma \in \text{Weld}(\Gamma)$ .

**Proof.** Given a quasicircle  $\Gamma \subset \mathbb{C}$  let  $F \in H(\Gamma)$ . Define  $G := F \circ \Phi$  and  $G_* := F \circ \Phi_* \circ \bar{h}$ . By the conformal invariance of the Dirichlet integral we have

$$(3.4) \quad \mathcal{D}[G] = \mathcal{D}_\Omega[F] \quad \text{and} \quad \mathcal{D}[G_*] = \mathcal{D}_{\Omega_*}[F] ,$$



and consequently, by (0.10),  $g := [g/\dagger]$ ,  $g_* := [g_*/\dagger] \in H$ , where  $g := \hat{\partial}G$  and  $g_* := \hat{\partial}G_*$ . Since both the functions  $g$  and  $g_*$  are continuous on  $\mathbf{T}$  and  $g = g_* \circ \gamma$ , we conclude from (0.12) that

$$(3.5) \quad g = B_\gamma(g_*) \quad \text{and} \quad g_* = B_\gamma^{-1}(g) .$$

From (3.4), (3.5) and (0.6) it follows that

$$\begin{aligned} \frac{\mathcal{D}_\Omega[F] - \mathcal{D}_{\Omega_*}[F]}{\mathcal{D}_\Omega[F] + \mathcal{D}_{\Omega_*}[F]} &= \frac{\mathcal{D}[G] - \mathcal{D}[G_*]}{\mathcal{D}[G] + \mathcal{D}[G_*]} = \frac{\|g\|_H^2 - \|g_*\|_H^2}{\|g\|_H^2 + \|g_*\|_H^2} \\ &= \frac{\|B_\gamma(g_*)\|_H^2 - \|g_*\|_H^2}{\|B_\gamma(g_*)\|_H^2 + \|g_*\|_H^2} \leq \frac{\|B_\gamma\|^2 - 1}{\|B_\gamma\|^2 + 1} , \end{aligned}$$

and similarly

$$\frac{\mathcal{D}_{\Omega_*}[F] - \mathcal{D}_\Omega[F]}{\mathcal{D}_\Omega[F] + \mathcal{D}_{\Omega_*}[F]} = \frac{\|B_\gamma^{-1}(g)\|_H^2 - \|g\|_H^2}{\|B_\gamma^{-1}(g)\|_H^2 + \|g\|_H^2} \leq \frac{\|B_\gamma^{-1}\|^2 - 1}{\|B_\gamma^{-1}\|^2 + 1} .$$

Combining the above inequalities we obtain

$$\frac{|\mathcal{D}_\Omega[F] - \mathcal{D}_{\Omega_*}[F]|}{\mathcal{D}_\Omega[F] + \mathcal{D}_{\Omega_*}[F]} \leq \frac{(\max\{\|B_\gamma\|, \|B_\gamma^{-1}\|\})^2 - 1}{(\max\{\|B_\gamma\|, \|B_\gamma^{-1}\|\})^2 + 1} ,$$

and consequently

$$(3.6) \quad \frac{1}{\lambda_s(\Gamma)} \leq \frac{(\max\{\|B_\gamma\|, \|B_\gamma^{-1}\|\})^2 - 1}{(\max\{\|B_\gamma\|, \|B_\gamma^{-1}\|\})^2 + 1} .$$

It remains to show the inverse inequality of (3.6). Fix  $g_* \in \text{Re}\mathcal{P}(\mathbf{T})$  and let  $g := g_* \circ \gamma$ . For  $z \in \hat{\mathbb{C}}$  define  $F(z) := \mathcal{P}[g] \circ \hat{\Phi}(z)$  if  $z \in \Omega$ ,  $F(z) := \mathcal{P}[g_*] \circ \bar{h} \circ \hat{\Phi}_*(z)$  if  $z \in \Omega_*$  and  $F(z) := \hat{\partial}(\mathcal{P}[g] \circ \hat{\Phi}(z))$  if  $z \in \Gamma$ . Since both the functions  $g$  and  $g_*$  are continuous on  $\mathbf{T}$ , we see that for every  $z \in \Gamma$

$$\hat{\partial}(\mathcal{P}[g] \circ \hat{\Phi})(z) = g \circ \hat{\partial}\hat{\Phi}(z) = g_* \circ \gamma \circ \hat{\partial}\hat{\Phi}(z) = g_* \circ \hat{\partial}\hat{\Phi}_*(z) = \hat{\partial}(\mathcal{P}[g_*] \circ \bar{h} \circ \hat{\Phi}_*)(z) .$$

Therefore the function  $F$  is continuous on  $\hat{\mathbb{C}}$ .

We can assume that  $\|g_*\|_H = 1$ , where  $g_* := [g_*/\dagger] \in H$ . By (0.12),  $g \doteq \mathcal{S}_0 B_\gamma(g_*)$ . The conformal invariance of the Dirichlet integral now shows, by (0.6), that

$$\mathcal{D}_\Omega[F] = \mathcal{D}[\mathcal{P}[g]] = 2\|B_\gamma(g_*)\|_H^2 < \infty$$

and

$$\mathcal{D}_{\Omega_*}[F] = \mathcal{D}[\mathcal{P}[g_*]] = 2\|g_*\|_H^2 = 2 .$$

Hence  $F \in H(\Gamma)$  and

$$(3.7) \quad \frac{\|B_\gamma(\mathbf{g}_*)\|_H^2 - 1}{\|B_\gamma(\mathbf{g}_*)\|_H^2 + 1} \leq \frac{|\mathcal{D}_\Omega[F] - \mathcal{D}_{\Omega_*}[F]|}{\mathcal{D}_\Omega[F] + \mathcal{D}_{\Omega_*}[F]} \leq \frac{1}{\lambda_s(\Gamma)}.$$

By (0.9),  $\|B_\gamma\| = \sup\{\|B_\gamma(\mathbf{f})\|_H : \mathcal{S}_0(\mathbf{f}) \in \text{Re } \mathbb{P}(\mathbf{T}) \text{ and } \|\mathbf{f}\|_H = 1\}$ . Then (3.7) leads to

$$\frac{\|B_\gamma\|^2 - 1}{\|B_\gamma\|^2 + 1} \leq \frac{1}{\lambda_s(\Gamma)}.$$

The same conclusion can be drawn for the inverse operator  $B_\gamma^{-1}$ , and so

$$\frac{(\max\{\|B_\gamma\|, \|B_\gamma^{-1}\|\})^2 - 1}{(\max\{\|B_\gamma\|, \|B_\gamma^{-1}\|\})^2 + 1} \leq \frac{1}{\lambda_s(\Gamma)}.$$

Combining this with (3.6) we obtain (3.3).  $\square$

Assume  $\Gamma \subset \mathbb{C}$  is a quasicircle. By (3.2) we see that  $\Psi$  is an extremal  $K$ -qc. reflection in  $\Gamma$  iff  $\psi := \bar{h} \circ \bar{\Phi}_* \circ \Psi \circ \Phi$  is an extremal  $K$ -qc. extension of  $\gamma$  to  $\Delta$ ,  $K \geq 1$ . Moreover, the complex dilatations of  $\psi$  and  $\Psi$  are related by the equality

$$(3.8) \quad \frac{\bar{\partial}\psi}{\partial\psi} = \frac{\overline{(\partial\Psi \circ \Phi)\Phi'}}{(\partial\Psi \circ \Phi)\bar{\Phi}'}$$
 a.e. on  $\Delta$ .

This observation is the key for the proof of the following

**Theorem 3.2.** *Given  $K > 1$  suppose that  $\Gamma \subset \mathbb{C}$  admits a regular  $K$ -qc. reflection  $\Psi$  with the complex dilatation*

$$(3.9) \quad \frac{\bar{\partial}\Psi}{\partial\Psi} = \frac{K + 1 \overline{G'}}{K - 1 G'} \quad \text{a.e. on } \Omega,$$

where  $G \in \dot{A}^2(\Omega)$  is a non-constant function. Then for each  $\gamma \in \text{Weld}(\Gamma)$

$$(3.10) \quad \lambda_s(\Gamma) = \frac{K + 1}{K - 1} = \frac{\|B_\gamma\|^2 + 1}{\|B_\gamma\|^2 - 1},$$

and  $\Psi$  is a unique extremal  $K$ -qc. reflection in  $\Gamma$ .

**Proof.** Given a non-constant function  $G \in \dot{A}^2(\Omega)$ , let  $F := G \circ \Phi - G \circ \Phi(0)$ . Then  $F \in \dot{A}(\Delta)$ ,  $\mathcal{D}[F] > 0$  and  $F' = (G' \circ \Phi)\Phi'$ . By (3.8) and (3.9) we have

$$(3.11) \quad \frac{\bar{\partial}\psi}{\partial\psi} = \frac{\overline{(\partial\Psi \circ \Phi)\Phi'}}{(\partial\Psi \circ \Phi)\bar{\Phi}'} = \frac{K - 1 \overline{(G' \circ \Phi)\Phi'}}{K + 1 (G' \circ \Phi)\bar{\Phi}'} = \frac{K - 1 \overline{F'}}{K + 1 F'} \quad \text{a.e. on } \Delta.$$

By (0.10),  $f := [\hat{\partial}_r \operatorname{Re} F / \ddot{=}] \in H$  and  $F = S_0(f)_\Delta$ .  
 By (0.11),  $4\|f\|_H^2 = \mathcal{D}[F] > 0$ . Theorem 1.3 now yields (1.29) for some  $g \in H$  satisfying  $\|g\|_H = 1$ . Therefore  $\|B_\gamma\| = \|B_\gamma^{-1}\|$  by Corollary 2.2. Hence (3.10) follows from (1.29) and Lemma 3.1. Moreover, from Theorem 1.3 and (3.11) we conclude that  $\psi$  is a unique extremal  $K$ -qc. extension of  $\gamma$  to  $\Delta$ , and hence that  $\Psi$  is a unique extremal  $K$ -qc. reflection in  $\Gamma$ .  $\square$

**Corollary 3.3.** *For every  $K > 1$  and every non-constant function  $G \in \dot{A}^2(\Omega)$  there exists a quasicircle  $\Gamma \subset \mathbb{C}$  which admits a unique extremal  $K$ -qc. reflection  $\Psi$  with the complex dilatation given by (3.9), and for each  $\gamma \in \operatorname{Weld}(\Gamma)$  the equality (3.10) holds.*

**Proof.** Fix a non-constant function  $G \in \dot{A}^2(\Omega)$ . Then the function  $F := G \circ \Phi - G \circ \Phi(0) \in \dot{A}(\Delta)$  is also non-constant. Following the proof of Corollary 1.4 we see that there exists  $\gamma \in \mathbb{Q}_T$  which admits a unique regular  $K(\gamma)$ -qc. Teichmüller extension  $\psi$  to  $\Delta$  with the complex dilatation given by (3.11), where  $K := K(\gamma)$ . It can be shown that  $\gamma$  is a welding homeomorphism of some quasicircle  $\Gamma \subset \mathbb{C}$ ; cf. e.g. [P1] or [V]. Then the mapping  $\Psi$ , given by (3.1), is a unique extremal  $K$ -qc. reflection and, by (3.8), its complex dilatation satisfies the equation (3.9). Then Theorem 3.2 shows that (3.10) holds for each  $\gamma \in \operatorname{Weld}(\Gamma)$ .  $\square$

**Theorem 3.4.** *Given  $K > 1$  suppose that  $\Gamma \subset \mathbb{C}$  admits a regular  $K$ -qc. reflection  $\Psi$ . If*

$$(3.12) \quad \lambda_s(\Gamma) = \frac{K + 1}{K - 1}$$

and if there exists a sequence  $G_n \in H(\Gamma)$ ,  $\mathcal{D}_\Omega[G_n] = 1$ ,  $n \in \mathbb{N}$ , such that

$$(3.13) \quad \frac{1}{\lambda_s(\Gamma)} = \lim_{n \rightarrow \infty} \frac{|\mathcal{D}_\Omega[G_n] - \mathcal{D}_{\Omega_*}[G_n]|}{\mathcal{D}_\Omega[G_n] + \mathcal{D}_{\Omega_*}[G_n]},$$

and that

$$(3.14) \quad \mathcal{D}_\Omega[G_n - G_m] \rightarrow 0 \quad \text{as } n, m \rightarrow \infty,$$

then the equation (3.9) holds for some non-constant function  $G \in \dot{A}^2(\Omega)$  and the equation (3.10) holds. In particular,  $\Psi$  is a unique extremal  $K$ -qc. reflection in  $\Gamma$ .

**Proof.** Let  $\gamma := \hat{\partial}\bar{\Phi}_* \circ \hat{\partial}\Phi \in \operatorname{Weld}(\Gamma)$ . For every  $n \in \mathbb{N}$ , set  $F_n := G_n \circ \Phi$ ,  $F_{n,*} := G_n \circ \Phi_* \circ \bar{h}$  and  $f_n := \hat{\partial}F_n$ . By the conformal invariance of the Dirichlet integral we have

$$(3.15) \quad \mathcal{D}[F_n] = \mathcal{D}_\Omega[G_n] = 1 < \infty \quad \text{and} \quad \mathcal{D}[F_{n,*}] = \mathcal{D}_{\Omega_*}[G_n] < \infty, \quad n \in \mathbb{N}.$$

Then for each  $n \in \mathbb{N}$ ,  $F_n \in \dot{H}^2(\Delta)$ , and consequently, by (0.10) and (0.11), we see that  $\mathbf{f}_n := [f_n / \doteq] \in \mathbf{H}$  and  $2\|\mathbf{f}_n\|_{\mathbf{H}}^2 = \mathcal{D}[F_n] = 1$ . Since each  $f_n \in C(\mathbb{T})$  and  $\hat{\partial}F_n = \hat{\partial}F_{n,*} \circ \gamma$ , we conclude from (0.12) and (0.6) that

$$(3.16) \quad 2\|B_\gamma^{-1}(\mathbf{f}_n)\|_{\mathbf{H}}^2 = \mathcal{D}[F_{n,*}], \quad n \in \mathbb{N}.$$

Applying the conformal invariance of the Dirichlet integral once again we see, by (3.14) and (0.6), that

$$2\|\mathbf{f}_n - \mathbf{f}_m\|_{\mathbf{H}}^2 = \mathcal{D}[F_n - F_m] = \mathcal{D}_\Omega[G_n - G_m] \rightarrow 0 \quad \text{as } n, m \rightarrow \infty.$$

Therefore there exists  $\mathbf{f} \in \mathbf{H}$  such that  $\|\mathbf{f} - \mathbf{f}_n\|_{\mathbf{H}} \rightarrow 0$  as  $n \rightarrow \infty$ , and hence  $\|\mathbf{f}\|_{\mathbf{H}} > 0$ . By this, (3.13), (3.15), (3.16) and by the continuity of the operator  $B_\gamma^{-1}$  we obtain

$$\begin{aligned} \frac{1}{\lambda_s(\Gamma)} &= \lim_{n \rightarrow \infty} \frac{|\mathcal{D}_\Omega[G_n] - \mathcal{D}_{\Omega_*}[G_n]|}{\mathcal{D}_\Omega[G_n] + \mathcal{D}_{\Omega_*}[G_n]} = \lim_{n \rightarrow \infty} \frac{|\mathcal{D}[F_n] - \mathcal{D}[F_{n,*}]|}{\mathcal{D}[F_n] + \mathcal{D}[F_{n,*}]} \\ &= \lim_{n \rightarrow \infty} \frac{|||B_\gamma^{-1}(\mathbf{f}_n)|||_{\mathbf{H}}^2 - |||\mathbf{f}_n|||_{\mathbf{H}}^2|}{|||B_\gamma^{-1}(\mathbf{f}_n)|||_{\mathbf{H}}^2 + |||\mathbf{f}_n|||_{\mathbf{H}}^2} = \frac{|||B_\gamma^{-1}(\mathbf{f})|||_{\mathbf{H}}^2 - |||\mathbf{f}|||_{\mathbf{H}}^2}{|||B_\gamma^{-1}(\mathbf{f})|||_{\mathbf{H}}^2 + |||\mathbf{f}|||_{\mathbf{H}}^2}. \end{aligned}$$

Hence by (3.12) we have

$$|||B_\gamma^{-1}(\mathbf{f})|||_{\mathbf{H}}^2 = K|||\mathbf{f}|||_{\mathbf{H}}^2 \quad \text{or} \quad |||B_\gamma^{-1}(\mathbf{f})|||_{\mathbf{H}}^2 = K^{-1}|||\mathbf{f}|||_{\mathbf{H}}^2,$$

and consequently

$$(3.17) \quad |||B_\gamma^{-1}(\mathbf{f})|||_{\mathbf{H}}^2 = K|||\mathbf{f}|||_{\mathbf{H}}^2 \quad \text{or} \quad |||B_\gamma(\mathbf{g})|||_{\mathbf{H}}^2 = K|||\mathbf{g}|||_{\mathbf{H}}^2,$$

where  $\mathbf{g} := B_\gamma^{-1}(\mathbf{f})$ . Set  $k := (K-1)(K+1)^{-1}$ . Suppose the second equality in (3.17) holds. Then  $K(\gamma) = K$  by Theorem 1.2 and (3.2). Theorem 1.3 now shows that  $\gamma$  admits a regular qc. Teichmüller extension  $\psi$  to  $\Delta$  with the complex dilatation (1.15) and  $\psi$  is uniquely extremal. Then the mapping  $\Psi$ , as given by (3.1), is a unique extremal  $K$ -qc. reflection and, by (3.8), its complex dilatation satisfies (3.9) with a non-constant function  $G := \mathcal{S}_0(\mathbf{f})_\Delta \circ \tilde{\Phi} \in \dot{A}^2(\Omega)$ . If the first equality in (3.17) holds, then by Corollary 2.2, the second equality in (3.17) holds for some  $\mathbf{g} \in \mathbf{H}$ ,  $|||\mathbf{g}||| > 0$ , and the rest of the proof runs as before.  $\square$

#### 4. Complementary remarks.

**Remark 4.1.** Theorem 1.3 states additionally that if  $\gamma \in \mathbb{Q}_\mathbb{T}$  admits a regular qc. Teichmüller extension  $\psi$  of to  $\Delta$  with the complex dilatation (1.14), then  $\psi$  is uniquely extremal, provided  $F$  is a square of an analytic

function which is square integrable on  $\Delta$ . In this way we have proved, by the way, a special case of Strebel's theorem; cf. [St1], [St2] and [LK, p. 153-154].

From Theorems 1.3 and 2.1 we obtain

**Remark 4.2.** Under the assumptions in the first part of Theorem 1.3 the following equality holds

$$(4.1) \quad \lambda_*(\gamma) = \frac{\|B_\gamma\|^2 + 1}{\|B_\gamma\|^2 - 1}.$$

From Remark 4.2 we get

**Remark 4.3.** Theorem 3.2 and Corollary 3.3 hold with the equality (3.10) replaced by

$$(4.2) \quad \lambda_s(\Gamma) = \lambda_*(\gamma) = \frac{K+1}{K-1}.$$

Every analytic Jordan curve  $\Gamma \subset \mathbb{C}$  is a quasicircle, which is clear e.g. from [LV, p. 97].

This can be also deduced from the inclusion  $\text{Weld}(\Gamma) \subset \mathbf{A}_T \subset \mathbf{Q}_T$  and (3.2).

Combining Kühnau's result [Kü1, Satz 5] with Theorem 3.2 yields

**Remark 4.4.** In case  $\Gamma$  is an analytic Jordan curve Theorem 3.2 is reduced to Kühnau's result [Kü2, p. 302].

The idea of using welding homeomorphisms in the study of topics covered by Section 3 appears in [Kü3], too.

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