

KEIKO FUJITA and MITSUO MORIMOTO

Conical Fourier transform of Hardy space of harmonic functions on the Lie ball

ABSTRACT. This paper is an extended version of a talk entitled "Hardy spaces of harmonic functions related to the complex sphere" and given at the 12-th Conference on Analytic Functions. The authors consider Hardy space of complex harmonic functions on the Lie ball with an inner product given by an integral on a part of the boundary of the Lie ball. They determine the image of the space under conical Fourier transformations.

1. Introduction. We denote \mathbb{R}^{n+1} by \mathbb{E} and \mathbb{C}^{n+1} by $\bar{\mathbb{E}}$. Let $z \cdot w = z_1 w_1 + \cdots + z_{n+1} w_{n+1}$, $z^2 = z \cdot z$, and $\|z\|^2 = z \cdot \bar{z}$. The open and the closed Lie balls of radius r with center at 0 are defined by

$$\bar{B}(r) = \{z \in \bar{\mathbb{E}} : L(z) < r\}, \quad 0 < r < \infty,$$

$$\bar{B}[r] = \{z \in \bar{\mathbb{E}} : L(z) \leq r\}, \quad 0 \leq r < \infty,$$

where $L(z) = \{\|z\|^2 + (\|z\|^4 - |z^2|^2)^{1/2}\}^{1/2}$ is the Lie norm. Note that $\bar{B}(\infty) = \bar{\mathbb{E}}$.

We denote by $\mathcal{O}(\bar{B}(r))$ the space of holomorphic functions on $\bar{B}(r)$ equipped with the topology of uniform convergence on compact sets and denote by

$\mathcal{O}(\bar{B}[r]) = \lim \text{ind}_{r' > r} \mathcal{O}(\bar{B}(r'))$ the space of germs of holomorphic functions on $\bar{B}[r]$. Put

$$\begin{aligned}\mathcal{O}_\Delta(\bar{B}(r)) &= \{f \in \mathcal{O}(\bar{B}(r)) : \Delta_z f(z) = 0\}, \\ \mathcal{O}_\Delta(\bar{B}[r]) &= \lim \text{ind}_{r' > r} \mathcal{O}_\Delta(\bar{B}(r')), \end{aligned}$$

where $\Delta_z = \partial^2/\partial z_1^2 + \partial^2/\partial z_2^2 + \cdots + \partial^2/\partial z_{n+1}^2$ is the complex Laplacian. We call an element of $\mathcal{O}_\Delta(\bar{B}(r))$ a complex harmonic function on $\bar{B}(r)$.

Let $n \geq 2$. We define the complex sphere with radius $\lambda \in \mathbb{C}$ by $\bar{S}_\lambda = \{z \in \bar{\mathbb{E}} : z^2 = \lambda^2\}$. If $\lambda = 0$, then \bar{S}_0 is called the complex light cone (or the complex isotropic cone). Put

$$\begin{aligned}\bar{S}_\lambda(r) &= \bar{S}_\lambda \cap \bar{B}(r), \quad |\lambda| < r \leq \infty, \\ \bar{S}_\lambda[r] &= \bar{S}_\lambda \cap \bar{B}[r], \quad |\lambda| \leq r < \infty, \\ \bar{S}_{\lambda,r} &= \partial \bar{S}_\lambda[r], \quad |\lambda| \leq r < \infty.\end{aligned}$$

If $|\lambda| < r$, then $\bar{S}_{\lambda,r}$ is a $(2n-1)$ -dimensional compact manifold on which the orthogonal group $SO(n+1)$ acts transitively. If $|\lambda| = r > 0$, then it reduces to the n -dimensional compact manifold $\bar{S}_{\lambda,r} = \bar{S}_\lambda[r] = \lambda S_1$, where $S_1 = \{x \in \mathbb{E} : x^2 = 1\}$ is the real unit sphere.

For $f, g \in \mathcal{O}_\Delta(\bar{B}[r])$ we put

$$(f, g)_{\bar{S}_{\lambda,r}} = \int_{\bar{S}_{\lambda,r}} f(z) \overline{g(z)} \dot{d}z,$$

where $\dot{d}z$ is the normalized invariant measure on $\bar{S}_{\lambda,r}$.

After some necessary preparation in Section 2 we show in Section 3 that $(f, g)_{\bar{S}_{\lambda,r}}$ is an inner product on $\mathcal{O}_\Delta(\bar{B}[r])$ and denote by $h_\lambda^2(\bar{B}(r))$ the completion of $\mathcal{O}_\Delta(\bar{B}[r])$ with respect to the inner product $(f, g)_{\bar{S}_{\lambda,r}}$. We can see that $h_\lambda^2(\bar{B}(r))$ is isomorphic to a Hardy space of harmonic functions on the Lie ball.

In Section 4, we define the conical Fourier transformation $\mathcal{F}_{\mu,r}^\Delta$ for $f \in \mathcal{O}_\Delta(\bar{B}(r))$, where μ is another complex number with $|\mu| \leq r$. Then the conical Fourier transform $\mathcal{F}_{\mu,r}^\Delta f$ is given by

$$\mathcal{F}_{\mu,r}^\Delta f(\zeta) = \int_{\bar{S}_{\mu,r}} \exp(sz \cdot \zeta) \overline{f(z/s)} \dot{d}z, \quad \zeta \in \bar{S}_0,$$

which does not depend on $s > 1$ sufficiently close to 1.

Then in Section 5, by introducing a Radon measure on \bar{S}_0 , we construct the inverse mapping $\mathcal{M}_{\mu,r}$ of the conical Fourier transformation $\mathcal{F}_{\mu,r}^\Delta$.

We also study a Hilbert space $\mathcal{E}^2(\bar{S}_0; \mu, \lambda, r)$ of entire functions on \bar{S}_0 which are square integrable with respect to the Radon measure.

Finally, in Section 6, we show that the image of $h_\lambda^2(\bar{B}(r))$ under the conical Fourier transformation $\mathcal{F}_{\mu,r}^\Delta$ is isomorphic to $\mathcal{E}^2(\bar{S}_0; \mu, \lambda, r)$ and we study a reproducing kernel for $\mathcal{E}^2(\bar{S}_0; \mu, \lambda, r)$.

2. Homogeneous harmonic polynomials. We denote by $\mathcal{P}_\Delta^k(\bar{\mathbb{E}})$ the space of k -homogeneous harmonic polynomials on $\bar{\mathbb{E}}$. The dimension of $\mathcal{P}_\Delta^k(\bar{\mathbb{E}})$ is given by

$$N(k, n) = \frac{(2k + n - 1)(k + n - 2)!}{k!(n - 1)!} = O(k^{n-1}).$$

Let $P_{k,n}(t)$ be the Legendre polynomial of degree k and of dimension $n + 1$. The coefficient $\gamma_{k,n}$ of the highest power of $P_{k,n}(t)$ is known as

$$\gamma_{k,n} = \frac{\Gamma(k + (n + 1)/2)2^k}{N(k, n)\Gamma((n + 1)/2)k!}$$

and $\overline{P_{k,n}(t)} = P_{k,n}(\bar{t})$. The harmonic extension $\bar{P}_{k,n}(z, w)$ of $P_{k,n}(z \cdot w)$ is defined by

$$\bar{P}_{k,n}(z, w) = \left(\sqrt{z^2}\right)^k \left(\sqrt{w^2}\right)^k P_{k,n}\left(\frac{z}{\sqrt{z^2}} \cdot \frac{w}{\sqrt{w^2}}\right).$$

Then $\bar{P}_{k,n}(z, w)$ is a symmetric k -homogeneous harmonic polynomial in z and in w . If $z^2 = 0$ or $w^2 = 0$, then $\bar{P}_{k,n}(z, w) = \gamma_{k,n}(z \cdot w)^k$.

Theorem 2.1 ([6, Theorem 5.2]). *Define the k -harmonic component f_k of $f \in \mathcal{O}_\Delta(\bar{B}(r))$ by*

$$f_k(z) = N(k, n) \int_{S_1} f(\rho\omega) P_{k,n}(z/\rho, \omega) d\omega, \quad 0 < \rho < r,$$

where $d\omega$ is the normalized invariant measure on S_1 .

Then $f_k \in \mathcal{P}_\Delta^k(\bar{\mathbb{E}})$ and $\sum_{k=0}^\infty f_k(z)$ converges to $f(z)$ in the topology of $\mathcal{O}_\Delta(\bar{B}(r))$. Moreover, we have

$$f = \sum_{k=0}^\infty f_k(z) \in \mathcal{O}_\Delta(\bar{B}(r)) \iff \limsup_{k \rightarrow \infty} \|f_k\|_{C(S_1)}^{1/k} \leq 1/r,$$

$$f = \sum_{k=0}^\infty f_k(z) \in \mathcal{O}_\Delta(\bar{B}[r]) \iff \limsup_{k \rightarrow \infty} \|f_k\|_{C(S_1)}^{1/k} < 1/r,$$

where $\|f_k\|_{C(S_1)} = \sup\{|f_k(x)| : x \in S_1\}$.

For $f, g \in \mathcal{O}_\Delta(\bar{B}[r])$ we define the sesquilinear form $(\cdot, \cdot)_{\tilde{S}_{\lambda,r}}$ by

$$(f, g)_{\tilde{S}_{\lambda,r}} = \int_{\tilde{S}_{\lambda,r}} f(z)\overline{g(z)} \dot{d}z, \quad |\lambda| \leq r,$$

where $\dot{d}z$ is the normalized invariant measure on $\tilde{S}_{\lambda,r}$.

For $f_k \in \mathcal{P}_\Delta^k(\bar{\mathbb{E}})$, $g_l \in \mathcal{P}_\Delta^l(\bar{\mathbb{E}})$, R. Wada [13] proved the relation

$$(1) \quad \int_{\tilde{S}_{\lambda,r}} f_k(z)\overline{g_l(z)} \dot{d}z = L_{k,\lambda,r} \int_{S_1} f_k(x)\overline{g_l(x)} \dot{d}x = \begin{cases} L_{k,\lambda,r} \int_{S_1} f_k(x)\overline{g_l(x)} \dot{d}x \\ 0, & (k \neq l), \end{cases}$$

where

$$L_{k,\lambda,r} \equiv \begin{cases} |\lambda|^2 P_{k,n} \left(\frac{1}{2} \left(\frac{r^2}{|\lambda|^2} + \frac{|\lambda|^2}{r^2} \right) \right), & \lambda \neq 0, \\ \frac{\gamma_{k,n}}{2^k} r^{2k}, & \lambda = 0. \end{cases}$$

Note that $L_{k,0,r} = \lim_{\lambda \rightarrow 0} L_{k,\lambda,r}$.

Lemma 2.2 ([3, Lemma 7.2]). $L_{k,\lambda,r}$ is a monotone increasing function in $|\lambda|$; that is, for $0 < |\lambda| < |\mu| < r$ and $k \neq 0$, we have

$$2^{-k} \gamma_{k,n} r^{2k} = L_{k,0,r} < L_{k,\lambda,r} < L_{k,\mu,r} < L_{k,r,r} = r^{2k}.$$

By Lemma 2.2, Theorem 2.1, and (1) we have

$$(f, g)_{\tilde{S}_{\lambda,r}} = \sum_{k=0}^{\infty} \int_{\tilde{S}_{\lambda,r}} f_k(z)\overline{g_k(z)} \dot{d}z = \sum_{k=0}^{\infty} \int_{S_1} f_k(x)\overline{g_k(x)} \dot{d}x < \infty.$$

Thus $(\cdot, \cdot)_{\tilde{S}_{\lambda,r}}$ is an inner product on $\mathcal{O}_\Delta(\bar{B}[r])$.

The sesquilinear form $(f, g)_{\tilde{S}_{\lambda,r}} = \sum_{k=0}^{\infty} \int_{\tilde{S}_{\lambda,r}} f_k(z)\overline{g_k(z)} \dot{d}z$ was defined for $f, g \in \mathcal{O}_\Delta(\bar{B}[r])$. However, by Theorem 2.1, for $f \in \mathcal{O}_\Delta(\bar{B}[r])$ and $g \in \mathcal{O}_\Delta(\bar{B}(r))$

$$\int_{\tilde{S}_{\lambda,r}} f(sz)\overline{g(z/s)} \dot{d}z = \sum_{k=0}^{\infty} \int_{\tilde{S}_{\lambda,r}} f_k(z)\overline{g_k(z)} \dot{d}z$$

is well-defined for $s > 1$ sufficiently close to 1 and does not depend on s . Sometime we set $s \cdot \int_{\tilde{S}_{\lambda,r}} f(z)\overline{g(z)} \dot{d}z = \int_{\tilde{S}_{\lambda,r}} f(sz)\overline{g(z/s)} \dot{d}z$ and call it the symbolic integral over $\tilde{S}_{\lambda,r}$. Thus we can extend $(f, g)_{\tilde{S}_{\lambda,r}}$ to a separately

continuous sesquilinear form on $\mathcal{O}_\Delta(\bar{B}[r]) \times \mathcal{O}_\Delta(\bar{B}(r))$ by the symbolic integral. Similarly we can extend $(f, g)_{\bar{S}_{\lambda, r}}$ to a separately continuous sesquilinear form on $\mathcal{O}_\Delta(\bar{B}(r)) \times \mathcal{O}_\Delta(\bar{B}[r])$. Therefore, we still have

$$\overline{(f, g)_{\bar{S}_{\lambda, r}}} = (g, f)_{\bar{S}_{\lambda, r}}$$

for $f \in \mathcal{O}_\Delta(\bar{B}[r])$ and $g \in \mathcal{O}_\Delta(\bar{B}(r))$ or for $f \in \mathcal{O}_\Delta(\bar{B}(r))$ and $g \in \mathcal{O}_\Delta(\bar{B}[r])$.

3. Hardy spaces of harmonic functions on the Lie ball. Let $|\lambda| \leq r$. We denote by $h_\lambda^2(\bar{B}(r))$ the completion of $\mathcal{O}_\Delta(\bar{B}[r])$ with respect to the inner product $(\cdot, \cdot)_{\bar{S}_{\lambda, r}}$. By the definition,

$$h_\lambda^2(\bar{B}(r)) = \left\{ f = \sum_{k=0}^{\infty} f_k : f_k \in \mathcal{P}_\Delta^k(\bar{\mathbb{E}}), \sum_{k=0}^{\infty} \|f_k\|_{\bar{S}_{\lambda, r}}^2 < \infty \right\}.$$

Further, as in the proof of Lemma 3.2 in [2], we can see that $h_\lambda^2(\bar{B}(r))$ is isomorphic to the Hardy space:

$$h_\lambda^2(\bar{B}(r)) = \left\{ f \in \mathcal{O}_\Delta(\bar{B}(r)) : \sup_{0 < t < 1} \int_{\bar{S}_{\lambda, r}} |f(tz)|^2 dz < \infty \right\}.$$

Proposition 3.1 ([4, Theorem 1.5]). *The Hardy space $h_\lambda^2(\bar{B}(r))$ is a Hilbert space being a direct sum of the finite dimensional subspaces $\mathcal{P}_\Delta^k(\bar{\mathbb{E}})$:*

$$h_\lambda^2(\bar{B}(r)) = \bigoplus_{k=0}^{\infty} \mathcal{P}_\Delta^k(\bar{\mathbb{E}}).$$

By using Lemma 2.2, we can prove the following

Theorem 3.2 ([4, Theorem 1.5]). *For $0 < |\lambda| < |\mu| < r$, we have*

$$\mathcal{O}_\Delta(\bar{B}[r]) \subset h_r^2(\bar{B}(r)) \subset h_\mu^2(\bar{B}(r)) \subset h_\lambda^2(\bar{B}(r)) \subset h_0^2(\bar{B}(r)) \subset \mathcal{O}_\Delta(\bar{B}(r)).$$

Now we consider the reproducing kernel. Since $|P_{k, n}(z, w)| \leq L(z)^k L(w)^k$ and $\lim_{k \rightarrow \infty} (L_{k, \lambda, r})^{1/k} = r^2$ for $|\lambda| \leq r$, the Poisson kernel

$$K_{\lambda, r}(z, w) = \sum_{k=0}^{\infty} \frac{N(k, n)}{L_{k, \lambda, r}} \bar{P}_{k, n}(z, \bar{w})$$

is a function on $\{(z, w) \in \bar{\mathbb{E}} \times \bar{\mathbb{E}} : L(z)L(w) < r^2\}$ and complex harmonic in z . It satisfies $K_{\lambda, r}(z, w) = K_{\lambda, r}(w, z)$. In particular, $K_{r, r}(z, w)$ is the classical Poisson kernel and the restriction of $K_{0, r}(z, w)$ on $\bar{S}_0 \times \bar{\mathbb{E}}$ is called the Cauchy kernel on \bar{S}_0 in [8]:

$$\begin{aligned} K_{r, r}(z, w) &= K_{1, 1}(z/r, w/r), \\ K_{1, 1}(z, \bar{w}) &= \frac{1 - z^2 \bar{w}^2}{(1 + z^2 \bar{w}^2 - 2z \cdot \bar{w})^{(n+1)/2}}, \\ K_{0, r}(z, w) &= K_{0, 1}(z/r, w/r), \\ K_{0, 1}(z, \bar{w})|_{\bar{S}_0 \times \bar{\mathbb{E}}} &= \frac{1 + 2zw}{(1 - 2zw)^n}. \end{aligned}$$

Using the Poisson kernel, we have the following integral representation for $f \in \mathcal{O}_\Delta(\bar{B}(r))$ (Theorem 3 in [7], see also [10] and [11]):

$$f(z) = s \int_{\bar{S}_{\lambda, r}} f(w) K_{\lambda, r}(z, w) \dot{d}w, \quad z \in \bar{B}(r).$$

For $f \in h_\lambda^2(\bar{B}(r))$ we have

Theorem 3.3 ([4, Theorem 1.5]). *The Poisson kernel $K_{\lambda, r}(z, w)$ is a reproducing kernel of $h_\lambda^2(\bar{B}(r))$ which means that for $f \in h_\lambda^2(\bar{B}(r))$ we have the following integral representation:*

$$f(z) = (f(w), K_{\lambda, r}(w, z))_{\bar{S}_{\lambda, r}} = \int_{\bar{S}_{\lambda, r}} f(w) K_{\lambda, r}(z, w) \dot{d}w, \quad z \in \bar{B}(r).$$

We denote by $L^2\mathcal{O}(\bar{S}_{\lambda, r})$ the closed subspace of the space of square integrable functions on $\bar{S}_{\lambda, r}$ generated by $\mathcal{H}^k(\bar{S}_{\lambda, r}) = \mathcal{P}_\Delta^k(\bar{\mathbb{E}})|_{\bar{S}_{\lambda, r}}$, $k = 0, 1, 2, \dots$. Then as a corollary of Theorem 3 in [7] and Theorem 3.3 we have

Corollary 3.4. *The restriction mapping α_λ gives the following linear topological isomorphisms:*

$$\begin{aligned} \alpha_\lambda : h_\lambda^2(\bar{B}(r)) &\xrightarrow{\sim} \mathcal{O}_\Delta(\bar{B}(r)), \\ \alpha_\lambda : \mathcal{O}_\Delta(\bar{\mathbb{E}}) &\xrightarrow{\sim} \mathcal{O}(\bar{S}_\lambda), \end{aligned}$$

where $\mathcal{O}(\bar{S}_\lambda)$ is the space of holomorphic functions on \bar{S}_λ equipped with the topology of uniform convergence on compact sets.

For related topics see [3].

4. Conical Fourier transformation. Let $\mathcal{O}'_{\Delta}(\bar{B}[r])$ (resp., $\mathcal{O}'_{\Delta}(\bar{B}(r))$) be the dual space of $\mathcal{O}_{\Delta}(\bar{B}[r])$ (resp., $\mathcal{O}_{\Delta}(\bar{B}(r))$). For $T \in \mathcal{O}'_{\Delta}(\bar{B}[r])$, we define the Poisson transformation $\mathcal{P}_{\mu,r}$ by $\mathcal{P}_{\mu,r} : T \mapsto \mathcal{P}_{\mu,r}T(w) = \langle T_{z,\mu,r}(z, w) \rangle$. Then we have the following

Theorem 4.1. *Let $0 < r < \infty$. The Poisson transformation establishes the following antilinear topological isomorphisms:*

$$\begin{aligned} \mathcal{P}_{\mu,r} : \mathcal{O}'_{\Delta}(\bar{B}[r]) &\xrightarrow{\sim} \mathcal{O}_{\Delta}(\bar{B}(r)), \\ \mathcal{P}_{\mu,r} : \mathcal{O}'_{\Delta}(\bar{B}(r)) &\xrightarrow{\sim} \mathcal{O}_{\Delta}(\bar{B}[r]). \end{aligned}$$

Further, for $T \in \mathcal{O}'_{\Delta}(\bar{B}(r))$ and $f \in \mathcal{O}_{\Delta}(\bar{B}(r))$, or for $T \in \mathcal{O}'_{\Delta}(\bar{B}[r])$ and $f \in \mathcal{O}_{\Delta}(\bar{B}[r])$, we have

$$(2) \quad \langle T, f \rangle = (f, \mathcal{P}_{\mu,r}T)_{\bar{S}_{\mu,r}}.$$

This can be proved similarly as Theorem 15 in [10].

Since $\Delta_z \exp(z \cdot \zeta) = 0$ for $\zeta \in \bar{S}_0$, we can define the conical Fourier-Borel transformation for $T \in \mathcal{O}'_{\Delta}(\bar{B}[r])$ by

$$(3) \quad \mathcal{F}_r^{\Delta} : T \mapsto \mathcal{F}_r^{\Delta}T(\zeta) = \langle T_z, \exp(z\zeta) \rangle, \quad \zeta \in \bar{S}_0.$$

Put

$$\text{Exp}(\bar{S}_0; (r)) = \left\{ f \in \mathcal{O}(\bar{S}_0) : \forall_{r' > r}, \exists_{C > 0} \text{ s.t. } |f(\zeta)| \leq C \exp(r' L^*(\zeta)) \right\},$$

$$\text{Exp}(\bar{S}_0; [r]) = \left\{ f \in \mathcal{O}(\bar{S}_0) : \forall_{r', r}, \exists_{C > 0} \text{ s.t. } |f(\zeta)| \leq C \exp(r' L^*(\zeta)) \right\},$$

where

$$L^*(\zeta) = \sup \{ |z\zeta| : L(z) \leq 1 \} = \{ (\|\zeta\|^2 + |\zeta|^2) / 2 \}^{1/2}$$

is the dual Lie norm. Then we have the following

Theorem 4.2. *The conical Fourier-Borel transformation \mathcal{F}_r^{Δ} gives the following linear topological isomorphisms:*

$$(i) \quad \mathcal{F}_r^{\Delta} : \mathcal{O}'_{\Delta}(\bar{B}[r]) \xrightarrow{\sim} \text{Exp}(\bar{S}_0; (r)), \quad 0 \leq r < \infty,$$

$$(ii) \quad \mathcal{F}_r^{\Delta} : \mathcal{O}'_{\Delta}(\bar{B}(r)) \xrightarrow{\sim} \text{Exp}(\bar{S}_0; [r]), \quad 0 < r \leq \infty.$$

(cf. Theorem 18 in [9]).

Now we define the conical Fourier transformation $\mathcal{F}_{\mu,r}^\Delta$ on $\mathcal{O}_\Delta(\bar{B}(r))$ by

$$\mathcal{F}_{\mu,r}^\Delta = \mathcal{F}_r^\Delta \circ (\mathcal{P}_{\mu,r})^{-1}.$$

Then for $f \in \mathcal{O}_\Delta(\bar{B}(r))$, by (2) and (3), we have

$$\mathcal{F}_{\mu,r}^\Delta f(\zeta) = (\exp(z\zeta), f(z))_{\tilde{S}_{\mu,r}}, \zeta \in \tilde{S}_0.$$

Lemma 4.3. For $f = \sum_{k=0}^\infty f_k$, $f \in \mathcal{O}_\Delta(\bar{B}(r))$ and $f_k \in \mathcal{P}_\Delta^k(\bar{\mathbb{E}})$, we have

$$(4) \quad \mathcal{F}_{\mu,r}^\Delta f(\zeta) = \sum_{k=0}^\infty \frac{L_{k,\mu,r}}{N(k,n)k!^{\gamma_{k,n}}} \bar{f}_k(\zeta),$$

where we put $\bar{f}_k(\zeta) = \overline{f_k(\bar{\zeta})}$ for $f_k \in \mathcal{H}^k(\tilde{S}_0) \equiv \mathcal{P}_\Delta^k(\bar{\mathbb{E}})|_{\tilde{S}_0}$.

Proof. We have

$$(5) \quad \exp(z\zeta) = \sum_{k=0}^\infty \frac{1}{k!^{\gamma_{k,n}}} \tilde{j}_k \left(i\sqrt{z^2} \sqrt{\zeta^2} \right) \bar{P}_{k,n}(z, \zeta),$$

where

$$\begin{aligned} \tilde{j}_k(t) &= \Gamma \left(k + \frac{n+1}{2} \right) (t/2)^{-(k+\frac{n-1}{2})} J_{k+\frac{n-1}{2}}(t) \\ &= \sum_{l=0}^\infty \frac{(-1)^l \Gamma(k + \frac{n+1}{2})}{\Gamma(k + \frac{n+1}{2} + l) l!} (t/2)^{2l} \end{aligned}$$

is the entire Bessel function (see [6]). Thus by Theorem 2.1 and (1), we get (4). \square

Theorems 4.1 and 4.2 imply the following

Theorem 4.4. Let $0 < r < \infty$. The conical Fourier transformation $\mathcal{F}_{\mu,r}^\Delta$ gives following antilinear topological isomorphisms:

$$\begin{aligned} \mathcal{F}_{\mu,r}^\Delta : \mathcal{O}_\Delta(\bar{B}(r)) &\xrightarrow{\sim} \text{Exp}(\tilde{S}_0; (r)), \\ \mathcal{F}_{\mu,r}^\Delta : \mathcal{O}_\Delta(\bar{B}[r]) &\xrightarrow{\sim} \text{Exp}(\tilde{S}_0; [r]). \end{aligned}$$

By Corollary 3.4 we may assume $f_k \in \mathcal{P}_\Delta^k(\bar{\mathbb{E}})$. Therefore by (4) and Theorem 2.1 we obtain the following proposition (see also [8, Thm. 12]):

Proposition 4.5. Let $f = \sum_{k=0}^{\infty} f_k \in \text{Exp}(\bar{S}_0; (r))$ and $f_k \in \mathcal{H}^k(\bar{S}_0)$. Then we have

$$f = \sum_{k=0}^{\infty} f_k \in \text{Exp}(\bar{S}_0; (r)) \iff \limsup_{k \rightarrow \infty} \|k! f_k\|_{C(\bar{S}_{0,1})}^{1/k} \leq r/2,$$

$$f = \sum_{k=0}^{\infty} f_k \in \text{Exp}(\bar{S}_0; [r]) \iff \limsup_{k \rightarrow \infty} \|k! f_k\|_{C(\bar{S}_{0,1})}^{1/k} < r/2,$$

where $\|f_k\|_{C(\bar{S}_{0,1})} = \sup\{|f_k(z)| : z \in \bar{S}_{0,1}\}$.

5. Radon measures on \bar{S}_0 . Let $\rho_{\mu,r}(t)$ be a function on $[0, \infty)$ satisfying

$$(6) \quad \int_0^{\infty} t^{2k} \rho_{\mu,r}(t) dt = \frac{(N(k, n)k!)^2 \gamma_{k,n} 2^k}{L_{k,\mu,r}}, \quad k = 0, 1, \dots,$$

and define the Radon measure $d\bar{S}_{0(\mu,r)}$ on \bar{S}_0 by

$$\int_{\bar{S}_0} f(\zeta) d\bar{S}_{0(\mu,r)}(\zeta) \equiv \int_0^{\infty} \int_{\bar{S}_{0,1}} F(t\zeta') d\zeta' \rho_{\mu,r}(t) dt.$$

Such a function $\rho_{\mu,r}$ does exist by a theorem of A. Duran [1]. In case of $|\mu| = r$, K. Ii [5] and R. Wada [12] constructed such a function $\rho_r(t)$ of exponential type $-\tau$ by means of the modified Bessel functions.

By Corollary 4.5 and $\lim_{k \rightarrow \infty} (L_{k,\lambda,r})^{1/k} = r^2$, for $F \in \text{Exp}(\bar{S}_0; [r])$ and $G \in \text{Exp}(\bar{S}_0; (r))$ (resp., $F \in \text{Exp}(\bar{S}_0; (r))$ and $G \in \text{Exp}(\bar{S}_0; [r])$) the integral

$$\int_{\bar{S}_0} F(\zeta) \overline{G(\zeta)} d\bar{S}_{0(\mu,r)}(\zeta)$$

is well-defined and it defines a separately continuous sesquilinear form on $\text{Exp}(\bar{S}_0; [r]) \times \text{Exp}(\bar{S}_0; (r))$ (resp., $\text{Exp}(\bar{S}_0; (r)) \times \text{Exp}(\bar{S}_0; [r])$). If $w \in \bar{S}_0$ and $z \in B(r)$, then the function $w \mapsto \exp(z \cdot w)$ belongs to $\text{Exp}(\bar{S}_0; [r])$. Therefore, for $F \in \text{Exp}(\bar{S}_0; (r))$ we can define $\mathcal{M}_{\mu,r}F(z)$ by

$$(7) \quad \mathcal{M}_{\mu,r}F(z) = \int_{\bar{S}_0} \exp(z\zeta) \overline{F(\zeta)} d\bar{S}_{0(\mu,r)}(\zeta), \quad z \in \bar{B}(r).$$

We denote by $\mathcal{M}_{\mu,r}$ the transformation $F \mapsto \mathcal{M}_{\mu,r}F$. By Theorem 2.1, (5) and (1) we have the following

Lemma 5.1. For $F = \sum_{k=0}^{\infty} F_k \in \text{Exp}(\tilde{S}_0; (r))$ and $F_k \in \mathcal{H}^k(\tilde{S}_0)$, we have

$$\mathcal{M}_{\mu,r}F(w) = \sum_{k=0}^{\infty} \frac{N(k,n)k!\gamma_{k,n}}{L_{k,\mu,r}} \overline{F_k}(w).$$

Theorem 5.2. The mapping $\mathcal{M}_{\mu,r}$ gives following antilinear topological isomorphisms and is inverse to the conical Fourier transformation $\mathcal{F}_{\mu,r}^{\Delta}$:

$$\begin{aligned} \mathcal{M}_{\mu,r} : \text{Exp}(\tilde{S}_0; (r)) &\xrightarrow{\sim} \mathcal{O}_{\Delta}(\tilde{B}(r)), \\ \mathcal{M}_{\mu,r} : \text{Exp}(\tilde{S}_0; [r]) &\xrightarrow{\sim} \mathcal{O}_{\Delta}(\tilde{B}[r]). \end{aligned}$$

Proof. By Lemmas 4.3 and 5.1 we have $\mathcal{M}_{\mu,r} \circ \mathcal{F}_{\mu,r}^{\Delta} f(z) = f(z)$ for $f \in \mathcal{O}_{\Delta}(\tilde{B}(r))$ and $\mathcal{F}_{\mu,r}^{\Delta} \circ \mathcal{M}_{\mu,r} = F(z)$ for $F \in \text{Exp}(\tilde{S}_0; (r))$. Thus $\mathcal{M}_{\mu,r}$ is bijective, whereas $\mathcal{M}_{\mu,r}$ and $\mathcal{F}_{\mu,r}^{\Delta}$ are inverse to each other. \square

For $\zeta, \xi \in \tilde{S}_0$ we put

$$E_{\mu,r}(\zeta, \xi) = \int_{\tilde{S}_0} \exp(z\zeta) \overline{\exp(z\xi)} dz.$$

Proposition 5.3. For $F \in \text{Exp}(\tilde{S}_0; (r))$ we have

$$(8) \quad F(\xi) = \int_{\tilde{S}_0} F(\zeta) \overline{E_{\mu,r}(\zeta, \xi)} d\tilde{S}_{0(\mu,r)}(\zeta).$$

Proof. Let $F = \sum_{k=0}^{\infty} F_k \in \text{Exp}(\tilde{S}_0; (r))$ and $F_k \in \mathcal{H}^k(\tilde{S}_0)$. Then

$$\begin{aligned} F(\xi) &= \mathcal{F}_{\mu,r}^{\Delta} \circ \mathcal{M}_{\mu,r} F(\xi) = \left(\exp(z\xi), \int_{\tilde{S}_0} \exp(z\zeta) \overline{F(\zeta)} d\tilde{S}_{0(\mu,r)}(\zeta) \right)_{\tilde{S}_{\mu,r}} \\ &= \int_{\tilde{S}_{\mu,r}} \exp(sz \cdot \xi) \overline{\int_{\tilde{S}_0} \exp(z/s \cdot \zeta) \overline{F(\zeta)} d\tilde{S}_{0(\mu,r)}(\zeta)} dz \\ &= \int_{\tilde{S}_0} \int_{\tilde{S}_{\mu,r}} \exp(z\xi) \overline{\exp(z\zeta)} dz F(\zeta) d\tilde{S}_{0(\mu,r)}(\zeta) \\ &= \int_{\tilde{S}_0} F(\zeta) \overline{E_{\mu,r}(\zeta, \xi)} d\tilde{S}_{0(\mu,r)}(\zeta), \end{aligned}$$

where $s > 1$ is sufficiently close to 1. \square

Now we employ the theorem of A. Duran ([1]) again, and there is a C^∞ function $\rho_{\mu,\lambda,\tau}(t)$ which satisfies

$$(9) \quad \int_0^\infty t^{2k} \rho_{\mu,\lambda,\tau}(t) dt = \frac{(N(k,n)k!)^2 \gamma_{k,n} 2^k L_{k,\lambda,\tau}}{(L_{k,\mu,\tau})^2}, \quad k = 0, 1, \dots$$

Define the Radon measure $d\bar{S}_{0(\mu,\lambda,\tau)}$ on \bar{S}_0 by

$$\int_{\bar{S}_0} F(\zeta) d\bar{S}_{0(\mu,\lambda,\tau)}(\zeta) = \int_0^\infty \left(\int_{\bar{S}_{0,1}} F(t\zeta') \dot{\zeta}' \right) \rho_{\mu,\lambda,\tau}(t) dt.$$

When $|\mu| = |\lambda|$, (9) reduces to (6), $\rho_{\mu,\lambda,\tau}(t)$ to $\rho_{\mu,\tau}(t)$ and $d\bar{S}_{0(\mu,\lambda,\tau)}$ to $d\bar{S}_{0(\mu,\tau)}$. Put

$$\mathcal{E}^2(\bar{S}_0; \mu, \lambda, \tau) = \left\{ F \in \mathcal{O}(\bar{S}_0) : \int_{\bar{S}_0} |F(\zeta)|^2 d\bar{S}_{0(\mu,\lambda,\tau)}(\zeta) < \infty \right\}.$$

When $|\mu| = |\lambda|$, we denote $\mathcal{E}^2(\bar{S}_0; \mu, \lambda, \tau)$ by $\mathcal{E}^2(\bar{S}_0; \mu, \tau)$.

Theorem 5.4. *The Hilbert space $\mathcal{E}^2(\bar{S}_0; \mu, \lambda, \tau)$ is a Hilbert space being a direct sum of the finite dimensional subspaces $\mathcal{H}^k(\bar{S}_0)$:*

$$\mathcal{E}^2(\bar{S}_0; \mu, \lambda, \tau) = \bigoplus_{k=0}^\infty \mathcal{H}^k(\bar{S}_0).$$

Proof. Let $F = \sum_{k=0}^\infty F_k(\zeta) \in \mathcal{E}^2(\bar{S}_0; \mu, \lambda, \tau)$ and $F_k \in \mathcal{H}^k(\bar{S}_0)$. By the definition of the Radon measure $d\bar{S}_{0(\mu,\lambda,\tau)}$, we have

$$(10) \quad \begin{aligned} \int_{\bar{S}_0} |F(\zeta)|^2 d\bar{S}_{0(\mu,\lambda,\tau)}(\zeta) &= \int_0^\infty \left(\int_{\bar{S}_{0,1}} F(t\zeta') \dot{\zeta}' \right) \rho_{\mu,\lambda,\tau}(t) dt \\ &= \int_0^\infty \sum_{k=0}^\infty t^{2k} (F_k, F_k)_{\bar{S}_{0,1}} \rho_{\mu,\lambda,\tau}(t) dt \\ &= \sum_{k=0}^\infty \frac{(N(k,n)k!)^2 \gamma_{k,n} 2^k L_{k,\lambda,\tau}}{(L_{k,\mu,\tau})^2} \|F_k\|_{\bar{S}_{0,1}}^2. \end{aligned}$$

This completes the proof. □

By (10) and Lemma 2.2 we have the following

Corollary 5.5. *If $|\mu_1| < |\mu_2| \leq r$, then*

$$\mathcal{E}^2(\tilde{S}_0; \mu_1, \lambda, r) \subset \mathcal{E}^2(\tilde{S}_0; \mu_2, \lambda, r).$$

If $|\lambda_1| < |\lambda_2| \leq r$, then

$$\mathcal{E}^2(\tilde{S}_0; \mu, \lambda_1, r) \supset \mathcal{E}^2(\tilde{S}_0; \mu, \lambda_2, r).$$

If $|\mu_1| = |\lambda_1|, |\mu_2| = |\lambda_2|$ and $|\mu_1| < |\mu_2|$, then

$$\mathcal{E}^2(\tilde{S}_0; \mu_1, r) \subset \mathcal{E}^2(\tilde{S}_0; \mu_2, r).$$

6. The $\mathcal{F}_{\mu,r}^\Delta$ image of $h_\lambda^2(\bar{B}(r))$. Now we consider the image of the Hardy space $h_\lambda^2(\bar{B}(r))$ under the conical Fourier transformation $\mathcal{F}_{\mu,r}^\Delta$.

Let $f \in h_\lambda^2(\bar{B}(r))$. Since $h_\lambda^2(\bar{B}(r)) \subset \mathcal{O}_\Delta(\bar{B}(r))$ and

$$(11) \quad \exp(z\zeta) = \sum_{k=0}^{\infty} \frac{(z\zeta)^k}{k!} = \sum_{k=0}^{\infty} \frac{\tilde{P}_{k,n}(z, \zeta)}{k! \gamma_{k,n}}$$

for $z \in \bar{\mathbb{E}}$ and $\zeta \in \tilde{S}_0$, we have

$$\begin{aligned} \mathcal{F}_{\mu,r}^\Delta f(\zeta) &= (\exp(z\zeta), f(z))_{\tilde{S}_{\mu,r}} \\ &= \sum_{k=0}^{\infty} \frac{1}{k! \gamma_{k,n}} \left(\tilde{P}_{k,n}(z, \zeta), f(z) \right)_{\tilde{S}_{\mu,r}} \\ (13) \quad &= \sum_{k=0}^{\infty} \frac{L_{k,\mu,r}}{k! \gamma_{k,n} N(k,n)} \bar{f}_k(\zeta) \\ &= \sum_{k=0}^{\infty} \frac{L_{k,\mu,r}}{k! \gamma_{k,n} L_{k,\lambda,r}} \left(\tilde{P}_{k,n}(z, \zeta), f(z) \right)_{\tilde{S}_{\mu,r}}. \end{aligned}$$

For $z \in \bar{\mathbb{E}}$ and $\zeta \in \tilde{S}_0$, put

$$(13) \quad e_\lambda^\mu(z, \zeta) = \sum_{k=0}^{\infty} \frac{L_{k,\mu,r}(z\zeta)^k}{L_{k,\lambda,r} k!} = \sum_{k=0}^{\infty} \frac{L_{k,\mu,r} \tilde{P}_{k,n}(z, \zeta)}{L_{k,\lambda,r} k! \gamma_{k,n}}.$$

If $\zeta \in \tilde{S}_0$ is fixed, then $e_\lambda^\mu(\cdot, \zeta)$ is a complex harmonic function on $\bar{\mathbb{E}}$. Hence $e_\lambda^\mu(\cdot, \zeta) \in h_\lambda^2(\bar{B}(r))$. Therefore by (12), for $f \in h_\lambda^2(\bar{B}(r))$, we have

$$(14) \quad \mathcal{F}_{\mu,r}^\Delta f(\zeta) = (\exp(z\zeta), f(z))_{\tilde{S}_{\mu,r}} = (e_\lambda^\mu(z, \zeta), f(z))_{\tilde{S}_{\lambda,r}}.$$

Theorem 6.1. *The conical Fourier transformation $\mathcal{F}_{\mu,r}^\Delta$ gives the antilinear unitary isomorphism:*

$$(15) \quad \mathcal{F}_{\mu,r}^\Delta : h_\lambda^2(\bar{B}(r)) \xrightarrow{\sim} \mathcal{E}^2(\bar{S}_0; \mu, \lambda, r).$$

Proof. Let $f = \sum_{k=0}^\infty f_k \in h_\lambda^2(\bar{B}(r))$, $f_k \in \mathcal{P}_\Delta^k(\bar{\mathbb{E}})$ and put $F(\zeta) = \mathcal{F}_{\mu,r}^\Delta f(\zeta)$. Then by Lemma 4.3, (9) and (1), we have

$$\begin{aligned} \int_{\bar{S}_0} |F(\zeta)|^2 d\bar{S}_{0(\mu,\lambda,r)}(\zeta) &= \int_{\bar{S}_0} |\mathcal{F}_{\mu,r}^\Delta f(\zeta)|^2 d\bar{S}_{0(\mu,\lambda,r)}(\zeta) \\ &= \sum_{k=0}^\infty \frac{(N(k,n)k!)^2 \gamma_{k,n} 2^k L_{k,\lambda,r}}{(L_{k,\mu,r})^2} \left(\frac{L_{k,\mu,r}}{N(k,n)k! \gamma_{k,n}} \right)^2 \|f_k\|_{\bar{S}_{0,\lambda}}^2 \\ &= \sum_{k=0}^\infty \frac{2^k L_{k,\lambda,r}}{\gamma_{k,n}} \|f_k\|_{\bar{S}_{0,1}}^2 = \sum_{k=0}^\infty \frac{L_{k,\lambda,r}}{L_{k,0,1}} \|f_k\|_{\bar{S}_{0,1}}^2 \\ &= \sum_{k=0}^\infty L_{k,\lambda,r} \|f_k\|_{\bar{S}_1}^2 = \sum_{k=0}^\infty \|f_k\|_{\bar{S}_{\lambda,r}}^2 \\ &= \|f\|_{\bar{S}_{\lambda,r}}^2 = \int_{\bar{S}_{\lambda,r}} |f(z)|^2 dz < \infty. \end{aligned}$$

□

Combining Theorem 6.1 with Theorems 3.2 and 4.4 , we obtain

Proposition 6.2. *Let $|\lambda| \leq r$ and $|\mu| \leq r$. Then we have*

$$\text{Exp}(\bar{S}_0; [r]) \subset \mathcal{E}^2(\bar{S}_0; \mu, \lambda, r) \subset \text{Exp}(\bar{S}_0; (r)).$$

Since $\mathcal{E}^2(\bar{S}_0; \mu, \lambda, r) \subset \text{Exp}(\bar{S}_0; (r))$, the inverse mapping of (15) is given by (7) and every $F \in \mathcal{E}^2(\bar{S}_0; \mu, \lambda, r)$ is represented by the integral formula (8). But we also have formulas corresponding to (7) and (8) in terms of the function $e_\lambda^\mu(z, \zeta)$ and the measure $d\bar{S}_{0(\mu,\lambda,r)}$:

Proposition 6.3. *Let $F \in \text{Exp}(\bar{S}_0; (r))$. Then we have*

$$\mathcal{M}_{\mu,r} F(z) = \int_{\bar{S}_0} e_\lambda^\mu(z, \zeta) \overline{F(\zeta)} d\bar{S}_{0(\lambda,\mu,r)}(\zeta).$$

Proof. By (7), (6), (9), (11) and (13), the statement easily follows. □

Theorem 6.4. *The function*

$$(i) \quad E_{\mu,\lambda,r}(\zeta, \xi) = (e_{\lambda}^{\mu}(z, \zeta), e_{\lambda}^{\mu}(z, \xi))_{\tilde{S}_{\lambda,r}}$$

is a reproducing kernel for $\mathcal{E}^2(\tilde{S}_0; \mu, \lambda, r)$; that is, for $F \in \mathcal{E}^2(\tilde{S}_0; \mu, \lambda, r)$ we have the following integral representation:

$$F(\xi) = \int_{\tilde{S}_0} F(\zeta) \overline{E_{\mu,\lambda,r}(\zeta, \xi)} d\tilde{S}_{0(\lambda,\mu,r)}(\zeta).$$

We have

$$(ii) \quad E_{\mu,\lambda,r}(\zeta, \xi) = \sum_{k=0}^{\infty} \frac{(L_{k,\mu,r})^2}{N(k,n)(k!\gamma_{k,n})^2 L_{k,\lambda,r}} \bar{P}_{k,n}(\zeta, \bar{\xi}).$$

The Poisson kernel $K_{\lambda,r}(z, w)$ can be given as follows:

$$(iii) \quad K_{\lambda,r}(z, w) = \int_{\tilde{S}_0} e_{\lambda}^{\mu}(z, \zeta) \overline{e_{\lambda}^{\mu}(w, \zeta)} d\tilde{S}_{0(\lambda,\mu,r)}(\zeta).$$

Proof. If we write down the formula $F(\xi) = \mathcal{F}_{\mu,r}^{\Delta} \circ \mathcal{M}_{\mu,r} F(\zeta)$ using the function $e_{\lambda}^{\mu}(z, \zeta)$ and (14), we get the reproducing formula (ii):

$$\begin{aligned} E_{\mu,\lambda,r}(\zeta, \xi) &= (e_{\lambda}^{\mu}(z, \zeta), e_{\lambda}^{\mu}(z, \xi))_{\tilde{S}_{\lambda,r}} \\ &= \sum_{k=0}^{\infty} \left(\frac{1}{k!\gamma_{k,n}} \frac{L_{k,\mu,r}}{L_{k,\lambda,r}} \right)^2 (\bar{P}_{k,n}(z, \zeta), \bar{P}_{k,n}(z, \xi))_{\tilde{S}_{\lambda,r}} \\ &= \sum_{k=0}^{\infty} \left(\frac{1}{k!\gamma_{k,n}} \frac{L_{k,\mu,r}}{L_{k,\lambda,r}} \right)^2 \frac{L_{k,\lambda,r}}{N(k,n)} \bar{P}_{k,n}(\zeta, \bar{\xi}) \\ &= \sum_{k=0}^{\infty} \frac{(L_{k,\mu,r})^2}{N(k,n)(k!\gamma_{k,n})^2 L_{k,\lambda,r}} \bar{P}_{k,n}(\zeta, \bar{\xi}). \end{aligned}$$

(iii)

$$\begin{aligned} &\int_{\tilde{S}_0} e_{\lambda}^{\mu}(z, \zeta) \overline{e_{\lambda}^{\mu}(w, \zeta)} d\tilde{S}_{0(\lambda,\mu,r)}(\zeta) \\ &= \sum_{k=0}^{\infty} \frac{(N(k,n)k!)^2 \gamma_{k,n} 2^k L_{k,\lambda,r}}{(L_{k,\mu,r})^2} \left(\frac{L_{k,\mu,r}}{k!\gamma_{k,n} L_{k,\lambda,r}} \right)^2 (\bar{P}_{k,n}(z, \zeta), \bar{P}_{k,n}(w, \zeta))_{\tilde{S}_{0,1}} \\ &= \sum_{k=0}^{\infty} \frac{N(k,n)^2 2^k}{\gamma_{k,n} L_{k,\lambda,r}} \frac{L_{k,0,1}}{N(k,n)} \bar{P}_{k,n}(z, \bar{w}) = \sum_{k=0}^{\infty} \frac{N(k,n)}{L_{k,\lambda,r}} \bar{P}_{k,n}(z, \bar{w}) = K_{\lambda,r}(z, w). \end{aligned}$$

□

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Faculty of Culture and Education received 23 November 1998
Saga University, Saga 840-8502, Japan

Department of Mathematics
International Christian University
3-10-2 Osawa, Mitaka-shi, Tokyo 181-8585, Japan

